

Задача №10 відповідь

$$\text{Дано рівняння руху } \ddot{x} - 2\gamma\dot{x} + \omega^2 = E_0 e^{i\omega t}$$

$$\ddot{x} - 2\gamma\dot{x} + \frac{\omega^2}{m} x = E_0 e^{i\omega t}$$

$$\ddot{x} - 2\gamma\dot{x} + \Omega^2 x = E_0 e^{i\omega t}$$

Всех частотах юсоджено підмніж, який створений усіма частинами зваж одного із них  $\Omega$  - відстань змінення та ефективна маса системи шукати розв'язок у вигляді  $x(t) = A e^{-i\omega t}$

$$A(-i\omega)^2 e^{-i\omega t} - 2\gamma(-i\omega)A e^{-i\omega t} + \Omega^2 A e^{-i\omega t} = E_0 e^{-i\omega t}$$

$$A(2\gamma i\omega + \Omega^2 - \omega^2) = E_0$$

$$A = \frac{E_0}{2\gamma i\omega + \Omega^2 - \omega^2}$$

$$\text{Дано: } \bar{F} = \pm E$$

$$nq\bar{i} = \pm E$$

$$\frac{nqE_0}{2\gamma i\omega + \Omega^2 - \omega^2} = \pm E_0$$

$$f(\omega) = \frac{nqE_0}{\Omega^2 - \omega^2 + 2\gamma i\omega}$$

$$(14) \text{ Braūmu } \frac{\partial}{\partial x} (\varphi \vec{A}) + \frac{\partial}{\partial y} (\varphi \vec{A}) + \frac{\partial}{\partial z} (\varphi \vec{A})$$

$$\frac{\partial}{\partial x} (\varphi \vec{A}) = \vec{A} \frac{\partial \varphi}{\partial x} + \varphi \frac{\partial \vec{A}}{\partial x} = \frac{\partial \varphi}{\partial x} \frac{x}{2} \vec{A} + \varphi \frac{d\vec{A}}{dz} \cdot \frac{x}{2}$$

$$\frac{\partial \varphi}{\partial x} = \frac{\partial \varphi}{\partial z} \frac{\partial z}{\partial x} = \frac{\partial \varphi}{\partial z} \cdot \frac{x}{2}$$

$$\begin{aligned}
 (\vec{A} \nabla) \varphi &= \vec{A} \left( \frac{\partial \varphi}{\partial x}, \frac{\partial \varphi}{\partial y}, \frac{\partial \varphi}{\partial z} \right) = \frac{\partial \varphi}{\partial x} \cdot \frac{x}{2} \vec{A} + \varphi \frac{x}{2} \frac{d\vec{A}}{dz} + \frac{\partial \varphi}{\partial z} \frac{y}{2} \vec{A} + \\
 &+ \varphi \frac{y}{2} \frac{d\vec{A}}{dz} + \frac{\partial \varphi}{\partial z} \frac{z}{2} \vec{A} + \varphi \frac{z}{2} \frac{d\vec{A}}{dz} = \\
 &= \frac{\partial \varphi}{\partial z} \frac{\vec{A}}{2} (x+y+z) + \frac{\varphi}{2} \frac{d\vec{A}}{dz} (x+y+z) = \\
 &= \frac{\varphi \vec{A} \cdot \vec{z}}{2} + \frac{\varphi}{2} \vec{A} \cdot \vec{z} = \frac{\vec{z}}{2} (\varphi \vec{A} + \varphi \vec{A})
 \end{aligned}$$

(15)  $\text{Stetigkeit von } \varphi \text{ auf } M$  wenn, wenn  
 $\varphi = \bar{a}[\vec{B} \vec{z}]$

$$\begin{aligned} L\vec{B}\vec{z} &= \vec{c} \quad \varphi = \bar{a}\vec{c}; \quad [d \times [\nabla_f \times \vec{f}]] = \\ &= \bar{\nabla}_f(\bar{a}\bar{f}) - \bar{f}(\bar{a}\bar{\nabla}_f) \Rightarrow \nabla_f(\bar{a}\bar{f}) = \\ &= [d \times [\nabla_f \times f]] + f(d\bar{\nabla}_f); \quad (1) \\ [f \times [d\nabla_a d]] &= \nabla_d(f d) - d(f\nabla_d) \Rightarrow \\ \Rightarrow \nabla_d(f d) &= [\bar{f} \times [\nabla_d d]] + d(f\nabla_d) \quad (2) \end{aligned}$$

$$\text{grad } \bar{f} \bar{d} = \nabla_f(f d) + \nabla_d(f d) = (1) + (2)$$

in generalisierung  $\bar{a} \rightarrow \bar{d}, \bar{c} \rightarrow \bar{f}, \bar{a} \cdot \text{rezip. Betr.}$

$$\begin{aligned} \text{grad } (\bar{a}\bar{c}) &= [\bar{a}[\nabla_c \bar{c}]] + [\bar{c}[\nabla_a \bar{a}]] + \\ &+ \bar{c}(\bar{a}\nabla_c) + \bar{a}(\bar{c}\nabla_a) \end{aligned}$$

speziell  $c = [\vec{B} \vec{z}]$

$$\begin{aligned} \text{grad } (\bar{a}[\vec{B} \vec{z}]) &= [\bar{a} \times [d \times [\vec{B} \times \vec{z}]]] + [[\vec{B} \vec{z}] \times \\ &\times [\nabla_d \bar{a}]] + [\vec{B} \times \vec{z}](d\nabla_d \vec{z}) + d([\vec{B} \times \vec{z}] \nabla_d) \end{aligned}$$

$$\text{wd } [\vec{B} \vec{z}] = 2B$$

$$\text{grad } (\bar{a}[\vec{B} \vec{z}]) = [\bar{a} \times \vec{B}] + [\vec{B} \times \vec{z}](a\nabla_d \vec{z})$$

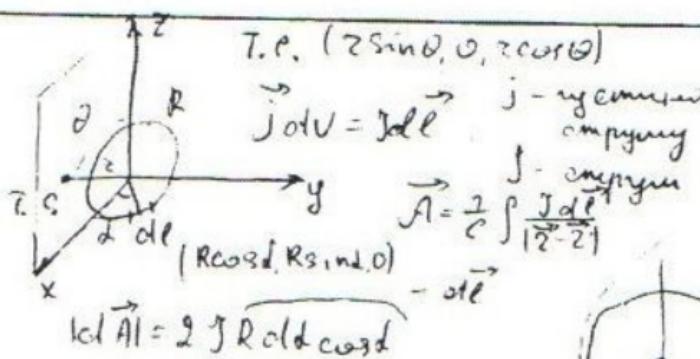
(16) Заряд є розподільний в атмосфері  $H$ , тає  
в нормальних умовах атмосфери є підігрів  
заряду  $\rho(z) = -\frac{\rho_0}{\pi a^3} \times \exp\left[-\frac{z^2}{a^2}\right]$ . Намагає  
чи,  $E_c$  - це, якщо ви зорогу, та маємо

$$\begin{aligned}
 E_c(z) &= \frac{4\pi}{z^2} \int_0^z z'^2 \left( -\frac{\rho_0}{\pi a^3} \right) \exp\left[-\frac{z'^2}{a^2}\right] dz' = \\
 &= -\frac{4\pi\rho_0}{\pi a^3 z^2} \int_0^z z'^2 e^{-\frac{z'^2}{a^2}} dz' = \left| \begin{array}{l} z' = u \\ dz' = 2z'dz \\ e^{-z'^2/a^2} dz = dV \end{array} \right| = \\
 &= -\frac{4\pi\rho_0}{a^3 z^2} \left[ -\frac{a^2}{2} e^{-\frac{z'^2}{a^2}} \Big|_0^z + 2\frac{a}{2} \int_0^z e^{-\frac{z'^2}{a^2}} z' dz \right] = \\
 &= -\frac{4\pi\rho_0}{a^3 z^2} \left[ -\frac{a}{2} z^2 e^{-\frac{z^2}{a^2}} + \frac{a}{2} \left( z^2 \left(-\frac{a}{2}\right) e^{-\frac{z^2}{a^2}} \right) \right] + \\
 &\quad + \frac{a}{2} \int_0^z e^{-\frac{z'^2}{a^2}} dz' \Big] = -\frac{4\pi\rho_0}{a^3 z^2} \left[ -\frac{a}{2} z^2 e^{-\frac{z^2}{a^2}} - \right. \\
 &\quad \left. - \frac{a^2}{2} z^2 e^{-\frac{z^2}{a^2}} - \frac{a^3}{4} e^{-\frac{z^2}{a^2}} + \frac{a^3}{4} \right] = \\
 &= -\frac{4\pi\rho_0}{a^3 z^2} \left[ \frac{\pi\rho_0}{4} e^{-\frac{z^2}{a^2}} \left( \frac{2}{a^2} + \frac{2}{a^2} + \frac{1}{2^2} \right) - \frac{\rho_0}{4 z^2} \right]
 \end{aligned}$$

$$E_c(z) = \rho_0 e^{-\frac{z^2}{a^2}} \left( \frac{2}{a^2} + \frac{1}{2^2} + \frac{2}{a^2} \right) - \frac{\rho_0}{2^2} = 4 \cdot$$

непаралельний заряд

(47) Three resistive load companies & galvanic  
source in the circuit



$$A_\phi = \frac{2\sqrt{R}}{c} \int_0^{\pi/2} \frac{\cos \alpha d\alpha}{\sqrt{2^2 + R^2 - 2R \cos \alpha \cos \beta}} = d\alpha$$

↑  $\rightarrow$   $d\alpha$

$$= \left| \begin{array}{l} \alpha = \pi/2 - \beta \\ d\alpha = -d\beta \quad \alpha = 0 \Rightarrow \beta = \pi/2 \\ \alpha = \pi/2 \Rightarrow \beta = 0 \\ \cos(\pi/2 - \beta) = \cos \beta = -1 + 2 \sin^2 \beta \end{array} \right|$$

↑  $\rightarrow$   $d\beta$

$$= \frac{2\sqrt{R}}{c} \int_0^{\pi/2} \frac{-(1 + 2 \sin^2 \beta - 1) d\beta}{\sqrt{2^2 + R^2 + 2R \sin \theta - 4R \sin \theta \sin^2 \beta}} =$$

$$= \frac{4\sqrt{R}}{c \sqrt{2^2 + R^2 + 2R \sin \theta}} \left[ -\frac{2}{k^2} \int_0^{\pi/2} \frac{1 - k^2 \sin^2 \beta - 1}{\sqrt{1 - k^2 \sin^2 \beta}} d\beta \right] -$$

$$- \left[ \frac{d\beta}{\sqrt{1 - k^2 \sin^2 \beta}} \right] \quad \text{E(k)} = \int_0^{\pi/2} \sqrt{1 - k^2 \sin^2 \beta} d\beta - \text{no.}$$

↑  $\rightarrow$   $d\beta$

↑  $\rightarrow$   $d\beta$

$$\ominus \frac{4\sqrt{R}}{c \sqrt{2^2 + R^2 + 2R \sin \theta}} \left[ -\frac{2}{k^2} E(k) + \left( \frac{2}{k^2} - 1 \right) K(k) \right]$$

$$\left\{ \vec{A}_\phi |_{\theta=0} = \frac{4\sqrt{R}}{c \sqrt{2^2 + R^2}} \left[ -\frac{2}{5} \frac{\pi}{2} + \left( \frac{2}{5} - 1 \right) \frac{\pi}{2} \right] \right\}$$

(18) Strain gradient ( $\vec{A}(r, \theta)$ )

$$z^2 = x^2 + y^2 + z^2 \quad \text{so } \frac{\partial z}{\partial x} = x/z$$

$$\frac{\partial \vec{A}(r, \theta)}{\partial x} = \frac{\partial}{\partial x} [A_x x + A_y y + A_z z] = \frac{\partial A_x}{\partial x} x + \frac{\partial A_y}{\partial y} y +$$

$$+ \frac{\partial A_z}{\partial x} z + A_x = \left( \frac{\partial A_i}{\partial x} = \frac{\partial A_i}{\partial r} \cdot \frac{\partial r}{\partial x} = A_i \frac{x}{r} \right) =$$

$$= \dot{A}_x \frac{x^2}{2} + \dot{A}_y \frac{y^2}{2} + \dot{A}_z \frac{z^2}{2} + A_x = \frac{x}{2} (\dot{A}_x x +$$

$$+ \dot{A}_y y + \dot{A}_z z) + A_x = \frac{x}{2} (\vec{A} \vec{z}) + A_x$$

$$\text{grad} (\vec{A}(r, \theta)) = \left\{ \frac{x}{2} (\vec{A} \vec{z}) + A_x; \frac{y}{2} (\vec{A} \vec{z}) + A_y; \right.$$

$$\left. \frac{z}{2} (\vec{A} \vec{z}) + A_z \right\} = \frac{(\vec{A} \vec{z})}{2} \cdot \{x, y, z\} + \{A_x, A_y, A_z\} =$$

$$= \frac{(\vec{A} \vec{z})}{2} \vec{z} + \vec{A}$$

straining grad  $\frac{\vec{P} \vec{z}}{2^3}$

$$x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

$$\vec{A} = \frac{\vec{P} \vec{z}}{2^3} = \frac{\vec{P}}{(x^2 + y^2 + z^2)^{1/2}} \{x, y, z\} = \frac{\vec{P}}{2^3} \{x, y, z\}$$

$$2^3 R = [2^2 \sin^2 \theta \cos^2 \phi + 2^2 \sin^2 \theta \sin^2 \phi + 2^2 \cos^2 \theta]^3/2 =$$

$$= 2^{6/2} = 2^3$$

$$\frac{\partial}{\partial z} \vec{A} = \vec{P} \left[ \frac{-2 \sin \theta \cos \phi}{2^3} + \frac{-2 \sin \theta \sin \phi}{2^3}, \frac{-2 \cos \theta}{2^3} \right]$$

$$\frac{\partial \vec{A}}{\partial \theta} = \frac{\vec{P}}{2^3} [2 \cos \theta \cos \phi + 2 \cos \theta \sin \phi - 2 \sin \theta]$$

$$\frac{\partial \vec{A}}{\partial \phi} = \frac{\vec{P}}{2^3} [-2 \sin \theta \sin \phi + 2 \sin \theta \cos \phi + 0]$$

$$\text{grad } \vec{A} = \vec{r}_z \frac{\partial \vec{A}}{\partial z} + \frac{\vec{e}_\theta}{2} \frac{\partial \vec{A}}{\partial \theta} + \frac{\vec{e}_\phi}{2 \sin \theta} \frac{\partial \vec{A}}{\partial \phi}$$

(19) Berechnen grad  $\vec{A}(z) \vec{B}(z)$

$$\frac{\partial \vec{A}(z) \vec{B}(z)}{\partial x} = \frac{\partial}{\partial x} [A_x B_x + A_y B_y + A_z B_z] =$$

$$= \frac{\partial}{\partial x} [A_i B_i] = \frac{\partial A_i}{\partial x} B_i + A_i \frac{\partial B_i}{\partial x} = \vec{A} \cdot \vec{B}; \frac{x}{2} +$$

$$+ \vec{B}_i A_i \frac{x}{2}$$

$$= \vec{A}_x B_x \frac{x}{2} + \vec{B}_x A_x \frac{x}{2} + \vec{A}_y B_y \frac{x}{2} + \vec{A}_y B_y \frac{x}{2} +$$

$$+ \vec{A}_z B_z \frac{x}{2} + \vec{A}_z B_z \frac{x}{2} = \frac{x}{2} \{ \vec{A}_x B_x + \vec{A}_y B_y + \vec{A}_z B_z +$$

$$+ \vec{B}_x A_x + \vec{B}_y A_y + \vec{B}_z A_z \} = \frac{x}{2} [\vec{A} \vec{B} + \vec{B} \vec{A}]$$

$$\text{grad } \vec{A}(z) \vec{B}(z) = \left\{ \frac{x}{2} [\vec{A} \vec{B} + \vec{B} \vec{A}], \frac{y}{2} [\vec{A} \vec{B}], \frac{z}{2} [\vec{A} \vec{B}] \right\}$$

Berechnen div  $\varphi(z) \vec{A}(z)$  =

$$= \varphi(z) \text{div } \vec{A}(z) + \vec{A}(z) \text{grad } \varphi(z) \oplus$$

$$\text{div } \vec{A}(z) = \vec{A}_x \frac{x}{2} + \vec{A}_y \frac{y}{2} + \vec{A}_z \frac{z}{2} = \frac{1}{2} (\vec{A} \vec{z})$$

$$\frac{\partial \vec{A}(z)}{\partial x} = \frac{\partial \vec{A}_x}{\partial z} \frac{\partial z}{\partial x} = \vec{A}_x \frac{x}{2}$$

$$\text{grad } \varphi(z) = \dot{\varphi} \left\{ \frac{x}{2}, \frac{y}{2}, \frac{z}{2} \right\} = \dot{\varphi} \frac{\vec{z}}{2}$$

$$\frac{\partial \varphi(z)}{\partial x} = \frac{\partial \varphi}{\partial z} \frac{\partial z}{\partial x} = \dot{\varphi} \frac{x}{2}$$

$$\oplus \frac{\varphi(z)}{2} (\vec{A} \cdot \vec{z}) + \vec{A}(z) \dot{\varphi} \frac{\vec{z}}{2} = \frac{\varphi(z)}{2} (\vec{A} \cdot \vec{z}) +$$

$$+ (\vec{A}(z) \vec{z}) \frac{\varphi}{2}$$

Berechnen rot  $\varphi(z) \vec{A}(z) = \varphi \text{rot } \vec{A} - \vec{A} \cdot \text{grad } \varphi \ominus$

$$\text{rot } \vec{A}(z) = \begin{vmatrix} \vec{e}_x & \vec{e}_y & \vec{e}_z \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \vec{A}_x & \vec{A}_y & \vec{A}_z \end{vmatrix} = \vec{e}_x \left[ \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right] -$$

$$- \vec{e}_y \left[ \frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z} \right] + \vec{e}_z \left[ \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right] =$$

$$= \vec{e}_x \left[ \vec{A}_z \frac{y}{2} - \vec{A}_y \frac{z}{2} \right] - \vec{e}_y \left[ -\vec{A}_x \frac{z}{2} + \vec{A}_z \frac{x}{2} \right] +$$

$$+ \vec{e}_z \left[ \vec{A}_y \frac{x}{2} - \vec{A}_x \frac{y}{2} \right] = \frac{1}{2} [\vec{z} \times \vec{A}]$$

$$\ominus \frac{\varphi}{2} [\vec{z} \times \vec{A}] + \frac{\varphi}{2} [\vec{z} \times \vec{A}]$$

(2c) Iloriumu nouluiu zilei emigratii



$$\frac{d\vec{B} \perp \vec{R}}{dI} = \frac{\mu_0 d\vec{e} \cdot \vec{R}}{R^3} = \frac{\mu_0 I}{R^2}$$

$$B = \frac{\mu_0 I R}{R^2} \Rightarrow B = \frac{\mu_0 I}{R} \text{ CGSM}$$

Hipoteza lui Ampere:  $\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{\text{int}}$



$$(3) \oint \vec{B} \cdot d\vec{s} = \mu_0 \int_R^R \mu_0 I dz$$

$$\text{c) } \text{z: } R \quad \vec{E} = \frac{e}{2\pi} \vec{z} \Rightarrow E = \frac{e}{2\pi} = \frac{\mu_0 I R^5}{5 z^2}$$

$$\varphi = \frac{e}{z} = \frac{\mu_0 I}{5\pi} R^5, \quad \varphi(R) = \frac{\mu_0 I}{5} R^4$$

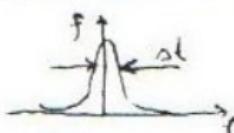
$$\text{d) } z \leq R \quad \oint \vec{B} \cdot d\vec{s} = \mu_0 \int_0^R \mu_0 I z^2 dz = \mu_0 \frac{4\pi}{5} \alpha z^5$$

$$D = \frac{4\pi}{5} \alpha z^3, \quad E = \frac{D}{\epsilon_0} = \frac{4\pi \alpha z^3}{5 \epsilon_0}$$

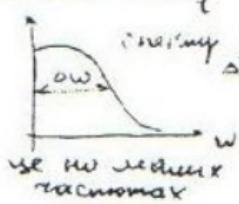
$$\varphi(z) - \varphi(R) = \int_R^z \frac{4\pi \alpha z^3}{5 \epsilon_0} dz = \frac{\pi \alpha}{5 \epsilon_0} (R^4 - z^4)$$

$$\varphi(z) = \frac{4\pi}{5} \alpha R^4 - \frac{\pi \alpha}{5 \epsilon_0} (R^4 - z^4)$$

(2) Частота  $\omega$  з-дана і є ціл. У цьому випадку обертанням об'єкта можна зберегти вимірювані засоби експерименту та зменшити узагор.



Тимчасові періоди є мінімальними узагор, а наявні зчлені  $\Rightarrow$  менше є мін. в мінімумі узагор  $\Rightarrow$  тащ-рів-л відхилення

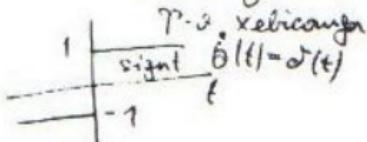
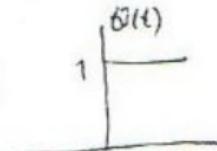


Ми поганоємо  $\omega = \omega_0 \sin(\theta + \omega t)$  макс. зчлені  $\Rightarrow$  добіг хвилі  $\rightarrow$  за це об'єкт  
засуб'єкт змін-ся та  
упоменевши

$$U_z = \begin{cases} U_0 \cos \theta, t < 0 \\ -U_0 \cos \theta, t > 0 \end{cases} \rightarrow U_0 \cos \theta \text{ sign}$$

$$\text{sign} = 2 \left( \theta - \frac{t}{2} \right)$$

$$V_z = 2 U_0 \cos \theta \left( \theta - \frac{t}{2} \right)$$



$$d = 2e, d_z = 2e \omega = \dot{\theta}_z e =$$

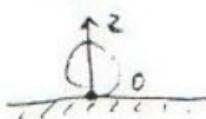
$$= -2e \left( U_0 \cos \theta \left( \theta - \frac{t}{2} \right) \right)' = -2e U_0 \cos \theta \delta(t)$$

$$(d_z)_w = - \int d_z e^{iwt} dt = -2e U_0 \cos \theta \int \delta(t) e^{iwt} dt =$$

$$= -2e U_0 \cos \theta$$

$$\frac{df}{dw} = \frac{2}{3\pi c^3} |\ddot{d}_w|^2 = \frac{2}{3\pi c^3} 4e^2 U_0^2 \cos^2 \theta =$$

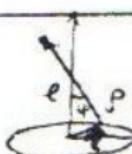
$$= \frac{8e^2}{3\pi c^3} U_0^2 \cos^2 \theta$$



Відображення мінімуму бунп.  
також.

(2)

$$\sigma = \text{const}$$



$$dS = 2\pi r d\theta d\phi$$

$$dE = \sigma^2 dS d\theta d\phi$$

$$dE = \frac{dE}{\rho^2} = \frac{\sigma^2 dS d\theta d\phi}{\rho^2}$$

$$dE \cos\phi = \frac{\sigma^2 dS d\theta d\phi}{\rho^2} \cos\phi$$

$$dE \cos\phi = 2\pi \sigma \frac{2\pi r^2 \cos\phi}{\rho^2}$$

$$r = \ell \sin\phi, \quad \rho = \ell / \cos\phi$$

$$dr = \frac{\ell}{\cos^2\phi} d\phi$$

$$dE \cos\phi = 2\pi \sigma \frac{\ell \sin\phi \ell d\phi \rho \cos\phi \cos^2\phi}{\cos^2\phi \ell^2} =$$

$$= 2\pi \sigma \sin\phi d\phi$$

$$E = 2\pi \sigma \left( 1 - \cos\phi \right) = 2\pi \sigma \left( 1 - \frac{\ell}{\sqrt{\ell^2 + R^2}} \right)$$

② Tipi ekrux yundax no-mg uwan. gunu.  
Cylinder. He zan. big cylinder noz. kwox.

$$\bar{M} = \frac{1}{2c} \int [ \vec{z}' \times \vec{j} (\vec{z}, t) ] dV, \quad \vec{z} = \vec{a} + \vec{z}'$$

$$\bar{M} = \frac{1}{2c} \int [ (\vec{a} + \vec{z}') \times \vec{j} (\vec{a} + \vec{z}') ] dV = \frac{1}{2c} \int [ \vec{z}' \times \vec{j} (\vec{a} + \vec{z}') ] dV' + \frac{1}{2c} \int [ \vec{a} \times \vec{j} (\vec{a} + \vec{z}') ] dV' = \bar{M}_0 + \frac{1}{2c} \int [ \vec{a} \times \vec{j} ] dV'$$

$$P = \frac{2}{3c^3} (\bar{M})^2 \quad P = \frac{2}{3c^3} (\bar{M}_0)^2$$

$$\bar{M} = \bar{M}_0 + \underbrace{\frac{1}{2c} \int [ \vec{a} \times \vec{j} ] dV'}_{=0} \Rightarrow \int \vec{j} dV = 0$$

$$\bar{d} = \int \rho \vec{z} dV, \quad \vec{d} = \int \vec{z} dV$$

$$\bar{d} = \int j \vec{z} dV, \quad \vec{d} = \int \vec{j} dV$$



$$\operatorname{div} \vec{j} = -\frac{\partial \rho}{\partial t} = -\dot{\rho} \quad d_x = \int j_x dV$$

$$(d \operatorname{div} \vec{j})_x = -\operatorname{div}(x \vec{j}) - \vec{j}_x$$

$$\begin{aligned} d_x &= \int (d \operatorname{div} \vec{j})_x dV = -\int \operatorname{div}(x_j) dV + \int j_x dV = \\ &= -\oint x_j \vec{n} dS + \int j_x dV \Rightarrow \vec{d} = \text{const} \end{aligned}$$

Знайти потенціал  $\varphi$   
 і напруженість  $E$   
 в ел. поля рівномірно  
 зарядженої прямолінійної лінії.

Нехай вона симетрична з віссю  $z$   
 маємо

$$\Delta \varphi = 4\pi \rho (2, \varphi, z) dz$$

$$\rho(2, \varphi, z) = 2\delta(z)$$

Оскільки, що розглянута задача не  
 залежить від  $z$ ,

$$\frac{1}{2} \frac{\partial}{\partial z} \left( 2 \frac{\partial \varphi}{\partial z} \right) = 2\delta(z)$$

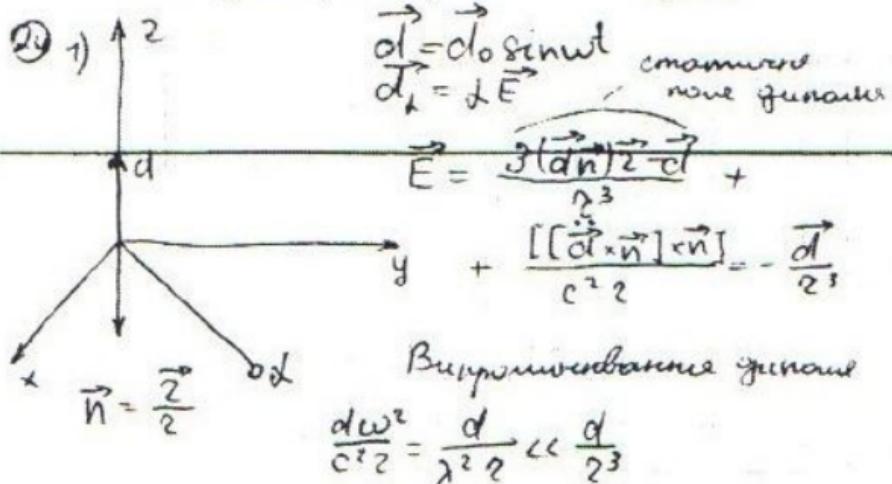
$$\frac{\partial}{\partial z} \left( 2 \frac{\partial \varphi}{\partial z} \right) = 2\delta(z)$$

$$\frac{\partial^2 \varphi}{\partial z^2} = \int 2\delta(z) dz = C_1$$

$$\frac{\partial \varphi}{\partial z} = \frac{C_1}{2}$$

$$\varphi = \int \frac{C_1}{2} dz = C_1 \ln|z| + C_2$$

$$E = -\operatorname{grad} \varphi = \frac{C_1}{z} \hat{z} \quad C_1 \neq 0$$



$\Rightarrow$  max repelling in equator

$$\Rightarrow \vec{E} = -\frac{\vec{\alpha}_\perp}{c^2 r^3}; \quad \vec{\alpha}_\perp = -\frac{\vec{d} \sin \omega t}{c^2 r^3}$$

$$\Rightarrow P = \frac{2}{3c^3} (\vec{\alpha}_\perp)^2 = \frac{2}{3c^3} \frac{d \vec{d}_0}{c^2 r^3} \omega d_0^2 \omega^4 \left(1 - \frac{d}{c^2 r^3}\right)^2 \times$$

$$x \sin^2 \omega t = \frac{2}{3c^3} \frac{d^2 d_0^2 \omega^4}{c^2 r^3} \sin^2 \omega t$$

2)  $\vec{e} = m \vec{v}_0$

$\vec{v}(t) = \vec{v}_0 \cos \theta \hat{x} + \vec{v}_0 \sin \theta \hat{y}$

$\vec{E} = \frac{P}{3c^3 d_0^2 \omega^4} \hat{z} = \frac{\vec{E}}{m}$

$\vec{\alpha} = \vec{e} \vec{v}(t), \quad P = \frac{2}{3c^3} (\vec{\alpha})^2$

$$\vec{F} = -\frac{2\vec{e}^2}{c^2} \Rightarrow \vec{v} = -\frac{2\vec{e}^2}{m c^2} \quad \vec{E} = \int p dl$$

$$P = \frac{2}{3c^3} \left( \ell \frac{2\vec{e}^2}{m c^2} \right)^2 = \frac{2\ell^2}{3c^3} \frac{e^6}{m^2 c^4}$$

$$E = \frac{2\ell^2 e^6}{3m^2 c^3} \int \frac{dl}{2^4(t)} = \frac{2\ell^2 e^6}{3m^2 c^3} \int \frac{dt}{(V_0^2 t^3 e^2)^2}$$

$$V_0 t = \sqrt{c^2 - \ell^2} \Rightarrow t = \sqrt{V_0^2 t^2 + \ell^2}$$

$$\cos \theta = \frac{\ell}{c}; \quad \tan \theta = \frac{V_0 t}{\ell}$$

(25)

Заряд распределен симмр. однозоне:  $\rho = \rho(z)$ . Равнное распределение заряда на симметрической оси выражено через  $\rho(z)$  момента

чтобы  $\varphi$  и  $E$  были.

$$E = 4\pi z^2 = 4\pi \int_{-\infty}^{\infty} \rho(z') 4\pi z'^2 dz' \Rightarrow E = \frac{4\pi}{2z} \int_{-\infty}^{\infty} \rho(z') z'^2 \times$$

$$z'^2 dz' = 4\pi \int_{-\infty}^{\infty} \frac{1}{z'^2} \int_{-\infty}^{\infty} \rho(z') z'^2 dz' dz' =$$

$$= \left| \frac{dV}{dz} = \frac{d\tilde{z}}{z^2} \right| = \frac{4\pi}{2} \int_{-\infty}^{\infty} \rho(z') z'^2 dz' \Big|_{z'} + 4\pi \int_{z}^{+\infty} \frac{1}{z^2} \times$$

$$\times \rho(\tilde{z}) \tilde{z}^2 d\tilde{z}$$

$$\varphi(z) = \frac{4\pi}{2} \int_{-\infty}^z \rho(z') z'^2 dz' + 4\pi \int_z^{\infty} \rho(\tilde{z}) \tilde{z}^2 d\tilde{z}$$

$$F(z) = \frac{4\pi}{2z} \int_0^z \rho(z') z'^2 dz'$$

Заряд распределен симметрически  
также // распределение симметрически симметрическое //  $\rho(z) = -\frac{C}{\pi z^3} \times \exp\left(-\frac{z^2}{a^2}\right)$ . Тогда  
 $\varphi, F$  — ли. волны, из-за  $\alpha$  (16).

26 Czynniki czyniące wóz ruchomy,  $y \propto \frac{e}{m} \text{ cm}$   
 Przyjmując zasada E. czer. teori. graw. uogólnioną  
 $\vec{d} = \sum_i e_i \vec{z}_i = \sum_i \frac{e_i}{m_i} (\vec{m}_i \vec{z}_i) = \frac{e}{m} \underbrace{\left( \sum_i (m_i \vec{z}_i) \right)}_M \vec{R}_g$   
 $(\sum m_i = M), \vec{d} = \frac{e}{m} M \vec{R}_g, \vec{d} = 0, \ddot{\vec{d}} = 0$   
 $p \sim (\vec{d})^2 = 0$

— II — II — teori. mechaniczno-graw. uogólnionych  
 $P_m = \frac{2}{3c^2} (\vec{M})^2$

$$\vec{M} = \frac{1}{2c} \sum_i [\vec{z}_i' \vec{j} (\vec{z}_i')]$$

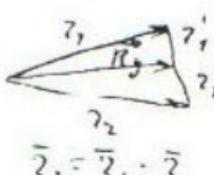
$$\vec{d} = \int \rho dV$$

$$\operatorname{div} \vec{j} = \frac{\partial \rho}{\partial t} = -\dot{\rho} \quad \vec{J} = \int j dV; \quad \vec{d} = \vec{j} \quad \text{wówczas masy...}$$

$$\vec{M} = \frac{1}{2c} \sum_i e_i [\vec{z}_i \times \frac{d\vec{z}_i}{dt}] = \frac{e}{2mc} \sum_i m [\vec{z}_i \times \vec{v}_i] = \\ = \frac{e}{2mc} \vec{L} = \text{const}$$

— II — II — 2 rozumygodz. — II — teori. mechaniczno-graw. uog.

$$\vec{M} = \frac{1}{2c} [e_1 [\vec{z}_1 \times \frac{d\vec{z}_1}{dt}] + e_2 [\vec{z}_2 \times \frac{d\vec{z}_2}{dt}]] = \\ = \frac{1}{2c} \left[ \frac{e_1 m_1}{m_1} [\vec{z}_1 \times \frac{d\vec{z}_1}{dt}] + \frac{e_2 m_2}{m_2} [\vec{z}_2 \times \frac{d\vec{z}_2}{dt}] \right] = \\ = \frac{1}{2c} \left[ \frac{e_1 \vec{L}_1}{m_1} + \frac{e_2 \vec{L}_2}{m_2} \right] \cdot \frac{1}{2c} \left[ \frac{e_1 m_1 \vec{L}_1 + e_2 m_2 \vec{L}_2}{m_1 m_2} \right]$$



$$\vec{R} = \frac{2}{m_1 + m_2} \vec{M} \\ \vec{z} = \vec{z}_1 - \vec{z}_2 \\ \vec{z}_1 = \vec{R} + \frac{m_2}{M} \vec{z}$$

$$\vec{R} = \frac{1}{M} (\vec{z}_1 m_1 + \vec{z}_2 m_2 - \vec{z} m_1 - \vec{z} m_2)$$



$$z = \frac{e}{\cos \theta}, \quad t = \frac{el \tan \theta}{V_0}, \quad dt = + \frac{l}{V_0} \frac{d\theta}{\cos^2 \theta}$$

$$E = \int \frac{z^2 e^6}{3m^2 c^3} \frac{e}{V_0} \frac{d\theta \cos^4 \theta}{\cos^2 \theta l^2} = \frac{z^2 e^6}{3m^2 e^3 l^3 V_0} \times$$

$$\times \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \cos^4 \theta d\theta = \frac{8}{2} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} (1 + \cos 2\theta) d\theta = \frac{8}{2} \times$$

$$\times \left[ \frac{\pi}{2} + \frac{\pi}{2} + \frac{1}{2} \sin^2 2\theta \right]_{-\frac{\pi}{2}}^{\frac{\pi}{2}} = \frac{8}{2} \left[ \frac{\pi}{2} (2) + \pi \right] =$$

$$= \frac{z^2 e^6 \pi}{3m^2 e^3 l^3 V_0}$$

$$T = \frac{mv^2}{2}, \quad \frac{E}{T} = \frac{z^2 e^6 \pi^2}{3m^2 e^3 l^3 V_0 m V_0^2} \quad (1)$$

$$V_0^3 \gg \frac{z^2 e^6}{l^3 m^3 e^3}; \quad V_0 \gg \frac{z^2 e^2}{cm l}$$

24 \*

(27) Розглянутий рухемоса навколо піраміди

$$\varphi(\bar{z}, t) = \frac{e}{R - \frac{(RV)}{c}} |_{t'} = t - \frac{|z - \bar{z}|}{c}$$

$$\vec{A}'(z, t) = \frac{e \bar{U}}{c(R - \frac{(RV)}{c})} |_{t'}$$

$$\vec{E} = \frac{c(e - \frac{U^2}{c^2})(R - \frac{\bar{U}}{c} R)}{(R - \frac{\bar{U} R}{c})^3} + \frac{e |R| ((\bar{R} - \frac{\bar{U}}{c} R) \bar{U}))}{c^2 (R - \frac{VR}{c})^2}$$

Розглянемо із загальн.

$$= \frac{1}{R} [\bar{R} E], \quad V=0$$

$$= \frac{e \bar{R}}{R^3}, \quad \text{тому } \frac{V}{c} \ll 1, \quad E \sim \frac{1}{R^2}; \quad H \sim \frac{1}{R^2}$$

$$B \sim \frac{1}{R^4} \quad [E \bar{H}] \sim \frac{1}{R^4}$$

Розглянемо із заг.

$$E = \frac{e R^2}{c^2 R^3} \sim \frac{1}{R}; \quad B \sim \frac{1}{R}$$

$$\vec{E} = -\operatorname{grad} \varphi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}, \quad \vec{B} = \omega \vec{A}$$

$$= \frac{m_1 + m_2}{m_1 m_2} [\bar{z}_1, (m_1 - m_2) - \bar{z}_2 m_2]$$

$$\bar{z}_1 = \bar{R} + \frac{m_2}{m} \bar{z}; \quad \bar{z}_2 = \bar{R} - \frac{m_1}{m} \bar{z}$$

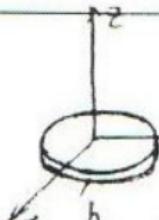
$$\bar{R} = 0; \quad \bar{z}_1 = \frac{m_2}{m} \bar{z}; \quad \bar{z}_2 = -\frac{m_1}{m} \bar{z};$$

$$\dot{\bar{z}}_{1,2} = \pm \frac{m_{1,2}}{m} \dot{\bar{z}}$$

$$\begin{aligned}\bar{M} &= \frac{1}{2c} \left( e_1 \left( \frac{m_2}{m} \right)^2 [\bar{z} \times \dot{\bar{z}}] + e_2 \left( \frac{m_1}{m} \right)^2 [\bar{z} \times \dot{\bar{z}}] \right) = \\ &= \frac{1}{2c} [\bar{z} \times \dot{\bar{z}}] \frac{1}{m^2} (e_1 m_2^2 + e_2 m_1^2) \sim \frac{1}{2c} \bar{N} \frac{1}{m^2} \times \\ &\times (e_1 m_2^2 + e_2 m_1^2) = \text{const}! \Rightarrow \bar{M} = 0 = \dot{\bar{M}}\end{aligned}$$

(26\*)

Q8. Задане квадратн. цил. с вис.  $h$ , вис. нас.  
q густина,  $\rho$ , ~~враща~~ h, піднамп. poz. zycomut. P.



$$\begin{cases} x = \rho \cos \varphi \\ y = \rho \sin \varphi \\ z = z \end{cases}$$

$$D_{xx} = \int_0^R \rho d\rho \int_0^{2\pi} d\varphi \int_{-\frac{h}{2}}^{\frac{h}{2}} dz x$$

$$\times (2\rho^2 \cos^2 \varphi - \rho^2 \sin^2 \varphi - z^2) =$$

$$= q \int_0^R \rho d\rho \int_0^{2\pi} (2\rho^2 \cos^2 \varphi h - \rho^2 \sin^2 \varphi h - \frac{h^3}{R}) =$$

$$= q \int_0^R \rho d\rho (2\rho^2 h - \overbrace{\rho^2 \pi h}^{0} - 2\pi \frac{h^3}{12}) = q \int_0^R (2\rho^2 \pi h - \frac{\pi h^3}{6}) d\rho :$$

$$q \left( 2 \frac{R^2 \pi h}{4} - \frac{\pi h^3 R^2}{12} \right) - \frac{d}{4} \left( R^2 - \frac{h^2}{3} \right)$$

$$D_{xx} - D_{yy} = \frac{1}{2} D_{zz} R$$

$$D_{xy} = D_{yx} = 0 = h \int_0^R \rho^3 d\rho \int_0^{2\pi} \sin 2\varphi d\varphi = 0$$

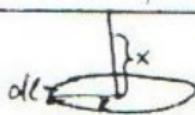
$$D_{xz} = D_{zx} = D_{yz} = D_{zy} = 0: \int_0^R \rho d\rho \int_{-\frac{h}{2}}^{\frac{h}{2}} dz \int_0^{2\pi} d\varphi \rho^2 \cos \varphi z$$

$$D_{zA} = \begin{pmatrix} \sum \ell_k (2x_k^2 - y_k^2 - z_k^2) & 3 \sum \ell_k x_k y_k & 3 \sum \ell_k x_k z_k \\ 3 \sum \ell_k x_k y_k & \sum \ell_k (2y_k^2 - x_k^2 - z_k^2) & 3 \sum \ell_k y_k z_k \\ 3 \sum \ell_k x_k z_k & 3 \sum \ell_k y_k z_k & \sum \ell_k (2z_k^2 - x_k^2 - y_k^2) \end{pmatrix}$$

$$D_{xx} = \sum \ell_k (3x_k^2 - y_k^2) - \sum \ell_k (2x_k^2 - y_k^2 - z_k^2)$$

(29)

Item E 8 yarangsi i na oci mindeg, surago  
 $\lambda = \lambda_0 \cos \varphi$



$$\Phi_B = \int_0^{\pi/2} \frac{R \lambda_0 \cos \varphi d\varphi}{R} =$$

$$= -\lambda_0 \varphi \sin \varphi \Big|_0^{\pi/2} = -\lambda_0 \frac{\pi}{2}$$

$$E = \frac{h \lambda}{R} \text{ ka oci } \varphi_0 = 4 \lambda_0 \int_0^{\pi/2} \frac{R \cos \varphi d\varphi}{\sqrt{R^2 + x^2}} =$$

$$= \frac{-4 \lambda_0}{\sqrt{R^2 + x^2}} \sin \varphi \Big|_0^{\pi/2} = \frac{-4 \lambda_0 R}{\sqrt{R^2 + x^2}}$$

$$x \rightarrow 0 \quad |E| = 4 \lambda_0$$

$$x \rightarrow \infty \quad |E| = \infty$$

$$|E| = -\text{grad } \varphi = \left| \frac{4 \lambda_0 R}{\sqrt{R^2 + x^2}} \right|_x = \frac{4 \lambda_0 R}{\sqrt{(R^2 + x^2)^2}}$$

$$|E| = 4 \lambda_0 \int_0^{\pi/2} \frac{R \cos \varphi d\varphi}{R^2 + x^2} + \frac{x}{\sqrt{R^2 + x^2}} = 4 \lambda_0 \frac{R}{\sqrt{R^2 + x^2}}$$

$$|E| = 4 \lambda_0 \sqrt{\frac{R \cos \varphi d\varphi}{R^2 + x^2}} \frac{x}{\sqrt{R^2 + x^2}} =$$

$$= 4 \lambda_0 \frac{R}{\sqrt{R^2 + x^2}}$$

$$E = \frac{4 \lambda_0 R \sqrt{R^2 + x^2}}{(x^2 + R^2)^{1/2}}$$

③ Енергоп. рух. премоніс. піднорматив. зона  
іменн. і енерг. наяв. близког. нормалізац.

$$\varphi(r,t) = \frac{e}{R - \frac{RV}{c}} \Big|_t, \quad A(r,t) = \frac{eVt}{c(R - \frac{RV}{c})} \Big|_t$$

З усього нормалізація зона. та. наяв.

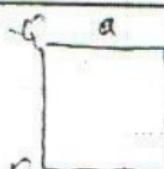
$$\bar{E} = \frac{e(c - V^2/c^2)/(R - \frac{V}{c}R)}{R - V} + \frac{e(R((R - \frac{V}{c}R) + t))}{c^2(R - \frac{V}{c}R)^3}$$

$$\bar{H} = \frac{1}{R}(\bar{R}\bar{E}), \quad V \ll c, \quad \bar{B} = \frac{e\bar{R}}{R^3}, \quad \text{також } \frac{V}{c} \ll 1$$

$$B = \frac{1}{R^2} \approx \sim \frac{1}{R^2}, \quad P \approx [\bar{B}\bar{H}] \sim \frac{1}{R^4}$$

Нормалізація з манетичною  $\phi = 0$  зона. наяв.

(34) Знайти співгучу по зажиг за репозиційною  $\zeta$  та залежністю від частоти



$$\text{D}_{\alpha\beta} = \sum (x_\alpha x_\beta \ddot{x}_i) q_i$$

$$\text{D}_{xx} = D_{yy} = Q \left( \frac{3a^2}{2} \cos^2 \omega t - \frac{a^2}{2} \right)$$

$$\text{D}_{yy} = -Q \left( \frac{3}{2} a^2 \sin^2 \omega t - \frac{a^2}{2} \right) + Q \left( \frac{3a^2}{2} \cos^2 \omega t - \frac{a^2}{2} \right) = 3Qa^2 \cos^2 \omega t$$

$$\text{D}_{yy} = +3a^2 Q \cos^2 \omega t$$

$$\text{D}_{zz} = -D_{xx} - D_{yy} = 0 \quad \text{cc}$$

$$x_1 = \frac{a}{\sqrt{2}} \cos \omega t, \quad y_1 = a \cos \omega t \sin \omega t$$

$$x_2 = \frac{a}{\sqrt{2}} \cos \left( \omega t - \frac{\pi}{4} \right), \quad y_2 = \frac{a}{\sqrt{2}} \sin \left( \omega t - \frac{\pi}{4} \right)$$

$$x_3 = \frac{a}{\sqrt{2}} \cos (\omega t - \frac{\pi}{2}), \quad y_3 = \frac{a}{\sqrt{2}} \sin (\omega t - \frac{\pi}{2})$$

$$x_4 = \frac{a}{\sqrt{2}} \cos \left( \omega t - \frac{3\pi}{2} \right), \quad y_4 = \frac{a}{\sqrt{2}} \sin \left( \omega t - \frac{3\pi}{2} \right)$$

$$y = \left( \frac{1}{180c^2} \text{D}_{\alpha\beta} \right) = \frac{1}{180c^2} (D_{xx}^2 + D_{yy}^2) = \frac{1}{180c^2} (2\omega)^6$$

$$\times (3a^2 Q \cos^2 \omega t) + (2\omega)^6 [ (3a^2 Q \cos^2 \omega t)^2 ] =$$

$$= \frac{2(7\omega)^6}{180c^2} 3^2 a^4 Q \cos^2 \omega t$$

$$\zeta = \frac{8}{5} \frac{\omega^6 a^4 Q}{c^5}$$

(35) Електронна модель роторного кривошипно-шатунного механизма  
 з вільним коливанням кривошипу розглянута виглядом  
роторної CLC як неподвижної K-силової рукої

$$C_1 = \gamma = c \quad \& \quad K' = A' = c \quad \varphi' = \ell / \sqrt{x'^2 + y'^2 + z'^2}$$

відповідно до вимірювань

$$\ell' = \frac{\varphi' + \frac{V}{c} A' c}{\sqrt{1 - \beta^2}} ; \quad A_{\perp} = \frac{A' + \beta A'}{\sqrt{1 - \beta^2}} ; \quad A_y = A'_y \quad A_x, A_z$$

$$y = \varphi' / \sqrt{1 - \beta^2} = \ell / [K' \sqrt{1 - \beta^2}] ; \quad A_{\perp} = c \beta \varphi$$

$$A_y = A_z = 0$$

Відповідно до цієї незадовільності

$$R'^2 = x'^2 + y'^2 + z'^2 = \frac{(x - U)^2}{1 - \beta^2} + y^2 + z^2 \Leftrightarrow$$

$$\Leftrightarrow \varphi = \frac{c}{R} \quad \text{де} \quad R = (x - U)^2 + (1 - \beta^2)(y^2 + z^2)$$

$$A = \varphi \beta$$

(36)  $\vec{B}$  - monce none dynm su emam. rukyo  
rot  $B=0$   $B[\vec{a} \times \vec{z}]$  - neperipisno

$$\nabla \times [\vec{a} \times \vec{z}], [\vec{a} \times \vec{z}] = \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ a_y & a_z & a_x \\ a_x & a_z & a_y \end{vmatrix} = i(a_y a_z - a_z a_y) + j(a_z a_x - a_x a_z) + k(a_x a_y - a_y a_x)$$

$$\nabla \times [\vec{a} \times \vec{z}] = \begin{vmatrix} i & j & k \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ (a_y a_z - a_z a_y) & (a_z a_x - a_x a_z) & (a_x a_y - a_y a_x) \end{vmatrix} =$$

$$= i(a_x + a_z) + j(a_z + a_x) + (a_x + a_z)k = 2i a_x + 2j a_y + 2k a_z = 2\pi - nre se su emamno$$

$$\text{rot rot } A = -\Delta A + \text{grad div } A = -\frac{1}{c} \frac{\partial^2 A}{\partial t^2}$$

Stream  $\vec{B}$  b eap i zebni

$$J = \epsilon_0 \sinh \left( \frac{r}{R} \right)^2$$

$$B 4\pi r^2 = 4\pi \int 4\pi r^2 J^2 dr \rho_0 \sinh^2 \left( \frac{r}{R} \right)^2 =$$

$$= \frac{4\pi r^2}{3} \int 4\pi J^2 \rho_0 \sinh^2 \left( \frac{r}{R} \right)^2 dr \left( \frac{r}{R} \right)^3 = \frac{16}{3} \pi^2 \rho_0 \times$$

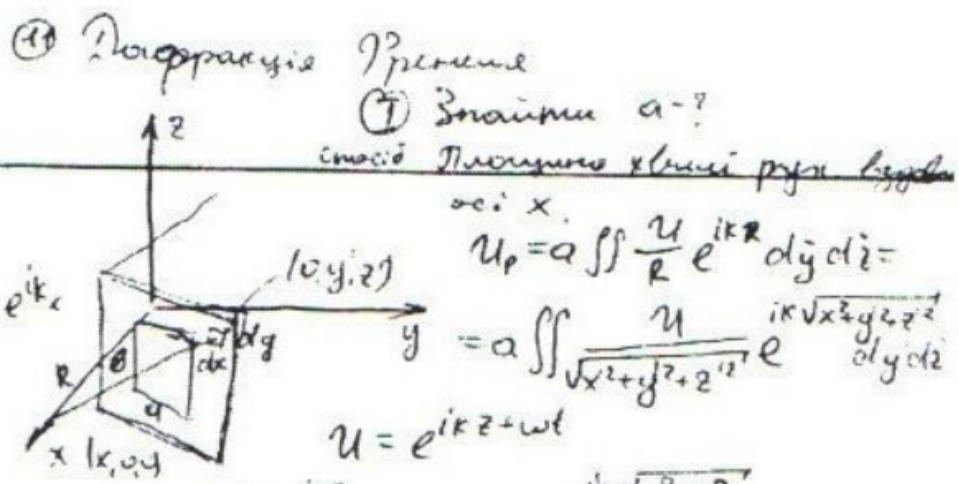
$$\times \left( \sinh^2 \left( \frac{r}{R} \right)^3 - 1 \right)$$

$$B_{\text{exp}} = \frac{4\pi \rho_0}{3r^2} \left( \sinh^2 \left( \frac{r}{R} \right)^3 - 1 \right) \quad \text{zebni} \quad r > R$$

$$B 4\pi r^2 = 4\pi \int 4\pi r^2 J^2 dr \left( \frac{r}{R} \right)^3 = \frac{16}{3} \pi^2 \rho_0 \times$$

$$\times \sinh^2 \left( \frac{r}{R} \right)^3 \Big|_0^R = \frac{16}{3} \pi^2 \rho_0 \left( \frac{r}{2} + \frac{1}{2e^2} - 1 \right)$$

$$B_{\text{zobni}} = \frac{4\pi \rho_0 \left( \frac{r}{2} + \frac{1}{2e^2} - 1 \right)}{3r^2}$$



$$U_p = a \int \frac{u e^{ikx}}{R} dz = a u \int \frac{e^{ik\sqrt{x^2 + z^2}}}{\sqrt{x^2 + z^2}} z dz d\Omega =$$

$$= 2\pi a u \int_0^\infty \frac{e^{ikx\sqrt{1 + \frac{z^2}{x^2}}}}{x} z dz = \frac{\pi a u}{x} \int_0^\infty e^{ikx(1 + \frac{z^2}{x^2})} z dz =$$

$$= \pi a \frac{u e^{ikx}}{x} \int_0^\infty e^{\frac{ikz^2}{x}} z dz = \pi a \frac{u e^{ikx}}{x} \frac{2x}{ik} e^{\frac{ikz^2}{x}} \Big|_0^\infty =$$

$$= -\frac{2\pi a}{ik} u e^{ikx} \Rightarrow \boxed{\alpha = \frac{k}{2\pi i}}$$

Young & unear long 0?

$\int f(x) dx$  - podzielne.  $\int f(k) e^{-ikx} dx$   $\rightarrow 0$  dla  $k \rightarrow \infty$ .

(II)  $R = \sqrt{x^2 + y^2 + z^2} = x \sqrt{1 + \frac{y^2}{x^2} + \frac{z^2}{x^2}} =$

$$x \times (1 + \frac{\frac{y^2}{x^2} + \frac{z^2}{x^2}}{2})$$

$$\dots \text{af } \frac{U_0 e^{ikx}}{R} dS = a \iint U_0 e^{ikx + \frac{1}{2} \frac{y^2}{x} + \frac{1}{2} \frac{z^2}{x}} dxdydz$$

$$= a U_0 e^{ikx} \int e^{ik \frac{y^2}{x}} d(\frac{y}{\sqrt{x}}) \int e^{ik \frac{z^2}{x}} d(\frac{z}{\sqrt{x}}) =$$

$$= -a U_0 e^{ikx} \sqrt{\frac{2\pi}{ik}} \sqrt{\frac{2\pi}{ik}} = \frac{2\pi i}{k} a U_0 e^{ikx} \Rightarrow$$

$$\Rightarrow \boxed{\alpha = \frac{k}{2\pi i}}$$

(2) Помітимо, що відповідно до розглянутого вище випадку, вектори  $\vec{d}_x$  та  $\vec{d}_y$  змінюються з часом за законом гармонічного коливання.

$$\vec{d}_x = d_0 \cos \omega t \vec{e}_x$$

$$\vec{d}_y = d_0 \sin \omega t \vec{e}_y$$

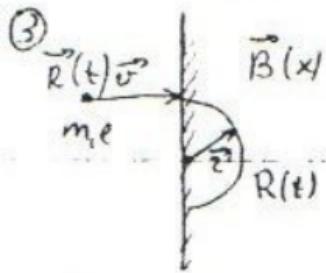
$$|\vec{d}|^2 = d_0^2 \theta^2 (\vec{e}_\theta \cdot \vec{e}_\theta)$$

$$\frac{dP}{d\Omega} = \frac{1}{4\pi c^3} \left( \omega^4 d_0^2 - \omega^4 d_0^2 \times \right. \\ \left. \times (\vec{n}_x \cos \omega t + \vec{n}_y \sin \omega t)^2 \right) \theta$$

$$\begin{cases} n_x = -\sin \theta \cos \varphi \\ n_y = \sin \theta \sin \varphi \end{cases} \Rightarrow \frac{1}{4\pi c^3} \omega^4 d_0^2 \left( 1 - \sin^2 \theta \times \right. \\ \left. \times (\cos \varphi \cos \omega t + \sin \varphi \sin \omega t)^2 \right)$$

$$= \frac{\omega^4 d_0^2}{4\pi c^3} \left( 1 - \sin^2 \theta \cos^2(\omega t - \varphi) \right)$$

$$\langle \frac{dP}{d\Omega} \rangle = \frac{\omega^4 d_0^2}{4\pi c^3} \left( 1 - \frac{1}{2} \sin^2 \theta \right) = \frac{\omega^4 d_0^2}{8\pi c^3} (1 + \cos^2 \theta)$$



Динамічні закони руху  
момента.

$$\vec{p} = \int \vec{r}' \rho(\vec{r}', t) d\Omega' = e \int \vec{r}' \times d\Omega' (\vec{r}' - R) d\Omega' = e \vec{R}(t)$$

$$\{ P = e \partial \gamma(\vec{r}', R(t)) \}$$

$$m \ddot{\vec{R}} = \vec{F}_n = \frac{e}{c} [\vec{U}, \vec{B}] ; [\vec{R}] : \frac{e}{mc} |[\vec{U}, \vec{B}]| = \frac{e}{mc}$$

$$\ddot{\vec{P}} = \frac{e^2}{mc} \vec{U} \vec{B} ; P_r = \frac{2}{3c^3} (\ddot{\vec{P}})^2 = \frac{2}{3c^3} \frac{e^4 U^2 B^2}{m^2 c^2} =$$

$$= \frac{2e^4}{3m^2 c^5} U^2 B^2 ; P = \frac{dE}{dt} \Rightarrow E_k = P \frac{T}{2} ; \omega = \frac{eB}{mc}$$

$$(a = \omega \nu)$$

$$\dot{E}_k = \frac{2e^4}{3m^2 c^5} U^2 B^2 \frac{\pi}{\omega} = \frac{2e^3}{3m^2 c^5} U^2 B^2 \frac{\pi T m e}{e B} = \\ = \frac{2\pi e^3 B U^2}{5mc^4}$$

Умов балансу високовимп. генератора  
не зумовлює ограничень на використання  
найменших значень зварювальних застосунків  
електро. застосувань.

$$E_k = \frac{m U^2}{2} \Rightarrow \frac{E_k}{E_B} \gg 1$$

$$\frac{E_k}{E_B} = \frac{m U^2 3m c^4}{4\pi e^3 B U^2} \gg 1 : B \ll \frac{3m^2 c^4}{4\pi e^3} \approx \frac{m^2 c^4}{e^3}$$

④ Томогузин ма мөн ма орци сунемпік  
гүйгү пайғызы R әркүншілік неберненделік үсмек-  
ма 3-мерзі  $\sigma = \delta_0 z^2$ ;  $\delta_0 = \text{const}$

$$\begin{aligned} d\ell &= \sqrt{z^2 + x^2 + y^2} = \sqrt{z^2 + r^2} \\ \varphi(z) &= \int \frac{\delta_0 z^2}{d\ell} = \\ &= \int_0^{R/2} \int_0^\pi \frac{\delta_0 z^2 r d\varphi}{\sqrt{z^2 + r^2}} dr = \\ &= 2\pi \delta_0 \int_0^{R/2} \frac{z^3 dr}{\sqrt{z^2 + r^2}} = \pi \delta_0 \int_0^{R/2} \frac{z^3 (z^2 + r^2)^{1/2}}{r} dr = |z^2 - y|^2 = \\ &= \pi \delta_0 \int_{R^2}^{z^2} \frac{y dy}{\sqrt{y^2 + z^2}} \quad (\text{E}) \\ \int \frac{y dy}{\sqrt{y^2 + z^2}} &= \int \sqrt{y^2 + z^2} dy = \int \frac{z^2 dy}{\sqrt{y^2 + z^2}} = \frac{z}{2} (y + z^2)^{1/2} - \\ &- \frac{z^2}{2} \sqrt{y^2 + z^2} \end{aligned}$$

$$\begin{aligned} (\text{E}) &\pi \delta_0 \left[ \frac{3}{2} (y + z^2)^{3/2} - \frac{z^2}{2} (y + z^2)^{1/2} \right] \Big|_{R^2} = \\ &= \pi \delta_0 \left[ \frac{3}{2} (y + z^2)^{3/2} - \frac{z^2}{2} (R^2 + z^2)^{1/2} - \frac{3}{2} z^3 + \frac{1}{2} z^5 \right] : \\ &= \frac{\pi \delta_0}{2} \left[ 3(R^2 + z^2)^{3/2} - z^2 (R^2 + z^2)^{1/2} - 2z^3 \right] \end{aligned}$$

$$\begin{aligned} E &= -\operatorname{grad} \varphi = -\frac{\partial \varphi}{\partial z} = -\frac{\pi \delta_0}{2} \left[ 3 \frac{3}{2} \sqrt{R^2 + z^2} z - \right. \\ &\quad \left. - 2z \sqrt{R^2 + z^2} - z^2 \frac{1}{2} (R^2 + z^2)^{-1/2} (2z - 6z^2) \right] = \\ &= -\frac{\pi \delta_0}{2} \left[ 7z \sqrt{R^2 + z^2} - \frac{z^3}{\sqrt{R^2 + z^2}} - 6z^3 \right] \end{aligned}$$

⑤ Знайти гарячійши момент зі сферами  
симетрії і ніч.

$$\begin{array}{c} \text{---} \\ -\lambda \quad +\lambda \\ \frac{\alpha}{2} \quad \alpha \end{array} \quad x \quad d = \int_{-\lambda}^{\lambda} -\lambda x dx + \int_{\lambda}^{\alpha} \lambda x dx = \\ = -\frac{\lambda}{2} \frac{\alpha^2}{4} + \frac{\lambda}{2} \left( \alpha^2 - \frac{\alpha^2}{4} \right) = \frac{\lambda \alpha^2}{4} \end{array}$$

$$= \frac{\lambda}{2} \left( \frac{3}{4} \alpha^2 - \frac{\alpha^2}{4} \right) = \frac{\lambda \alpha^2}{4}$$

$$\begin{array}{c} \text{---} \\ -\lambda \quad +\lambda \\ \frac{\alpha}{2} \quad \frac{\alpha}{2} \end{array} \quad x \quad d = - \int_{-\frac{\alpha}{2}}^{\frac{\alpha}{2}} \lambda x dx + \int_{\frac{\alpha}{2}}^{\frac{\alpha}{2}} \lambda x dx = \\ = \int_0^{\frac{\alpha}{2}} \lambda x dx + \int_0^{\frac{\alpha}{2}} \lambda x dx = \frac{\lambda}{2} \frac{\alpha^2}{4} + \frac{\lambda}{2} \frac{\alpha^2}{4} = \frac{\lambda \alpha^2}{4} \end{array}$$

Висновок: моменти двох сфер  
зі спільною осью симетрії

$$\begin{array}{l} dI = \frac{5 \pi dR^2 dz}{\sqrt{z^2 + R^2}} \\ I = 5 \int_0^R dz \int_0^{2\pi} \frac{2 \rho dz}{\sqrt{z^2 + R^2}} = \\ = 20 \pi \sqrt{z^2 + R^2} \Big|_0^R = 20 \pi (\sqrt{R^2 + z^2} - 1) \end{array}$$

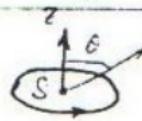
Знайти позначки відповідних зон

$$\begin{array}{c} e \\ a \end{array} \quad E(x) = 2 \frac{e}{a^2 + x^2} \cos(\theta) = \frac{2e}{a^2 + x^2} \times \\ \times \frac{x}{\sqrt{a^2 + x^2}} = \frac{2ex}{(a^2 + x^2)^{3/2}}$$

$$E = 4\pi G \Rightarrow$$

$$\langle E \rangle = \frac{E}{4\pi} = \frac{ea}{2\pi(a^2 + x^2)^{3/2}}$$

⑥ Einheitsweise notwendig zu erproben,  
 $\vec{p} = ?$   $\frac{d\vec{P}}{dt} = ?$



$$J = J_0 \cos \omega t$$

$$\vec{M} = \frac{JS}{C}, \vec{M} = \frac{1}{2C} \int [\vec{r} \cdot \vec{j}(\vec{r})] dV$$

$$\vec{M} = \frac{J_0 \cos \omega t S \hat{e}_z}{C}, \quad \ddot{\vec{M}} = - \frac{J_0 \omega^2 \cos \omega t S \hat{e}_z}{C}$$

$$\frac{dP_m}{d\Omega} = \frac{[\vec{M} \times \vec{n}]^2}{4\pi C^3} = \frac{1}{4\pi C^5} S^2 J_0^2 \omega^4 \cos^2 \omega t [\hat{e}_z \times \vec{n}]^2 = \\ = \frac{1}{4\pi C^5} S^2 J_0^2 \omega^4 \cos^2 \omega t \sin^2 \theta$$

$$\langle \frac{dP_m}{d\Omega} \rangle = \frac{1}{2} \left. \frac{dP_m}{d\Omega} \right|_{t=0} = \frac{J_0^2 S^2}{8\pi C^5} \omega^4$$

$$P_m = \frac{2}{3C^3} (\dot{\vec{M}})^2 = \frac{2}{3C^3} \frac{\omega^4 J_0^2 \cos^2 \omega t \cdot S^2}{C^2}$$

$$\vec{P}_m \Big|_{t=0} = \frac{2}{3C^5} \omega^4 J_0^2 S^2$$

(7) Trennung  $\propto$  unabhg., B - ?

Int:

$$dB = \frac{J}{c} \frac{\sin \alpha}{2^2} = \frac{J}{c} \frac{\sin \alpha}{2^2} =$$
$$= \frac{J}{c^2} \frac{d\cos \alpha}{2} = - \frac{J}{c^2} dz =$$
$$= \frac{J}{cR} \cos \alpha d\alpha$$
$$z = \frac{R}{\cos \alpha}, B = \int_{-\pi/2}^{\pi} \frac{J}{cR} \cos \alpha d\alpha = \frac{2J}{cR}$$

$\infty$  mehrere Anteile gebündelt  $\rightarrow$  superpos.

q. Maximum of max E  $\propto$  m. unabhg.

$$\ell_0 = \ell_1 + \ell_2, dE = \frac{d\ell}{2^2}$$
$$\ell = \ell_1 \ell_0$$
$$\ell = z_0 \tan \alpha, d\ell = \frac{z_0}{\cos^2 \alpha} d\alpha$$
$$z = \frac{z_0}{\cos \alpha}$$
$$dE = \frac{\ell_1 z_0 \alpha_1 \cos \alpha}{\cos^2 \alpha + z_0^2} =$$
$$= \frac{\ell_1}{z_0} d\alpha$$

$$dE_2 = dE \cos \alpha = \frac{\ell_1 \cos \alpha}{z_0} d\alpha$$

$$E_2 = \int_{\alpha_1}^{\alpha_2} \frac{\ell_1 \cos \alpha d\alpha}{z_0} = \frac{\ell_1}{z_0} (\sin \alpha_1 + \sin \alpha_2) =$$
$$= \frac{\ell_1}{z_0} \left( \frac{\ell_1}{\sqrt{\ell_1^2 + z_0^2}} + \frac{\ell_2}{\sqrt{\ell_2^2 + z_0^2}} \right) \propto$$

$$dF_d = -\frac{\ell_1}{z_0} \sin \alpha d\alpha \Rightarrow \frac{\ell_1}{z_0} (\cos \alpha_1 - \cos \alpha_2) = \dots$$

Statisch auf einer konzentrischen  $\frac{\pi}{2}$ ,  $\ell_1 = 0$

$$E_2 = \frac{\ell_1}{z_0}, E_d = -\frac{\ell_1}{z_0} \Rightarrow E = \frac{\sqrt{2} \ell_1}{z_0}$$

⑧ Напірна труба з певними вимірювальними



$$2) P = \alpha \cdot \epsilon \quad a) z \leq R$$

$$(1) P = \text{const} \quad \Rightarrow \quad 2\pi r^2 E = 4\pi r^2 \epsilon P$$

$$E = 2\pi r^2 P$$

$$\delta) z > R \quad 2\pi r^2 E = 4\pi R^2 \epsilon P, \quad E = \frac{2\epsilon P}{r}$$

$\ell_1$  - з-як має зменшувати гублення

$$g) \psi(z) - \psi(R) = \int_R^z E dz = \pi P (R^2 - z^2)$$

$$\delta) \psi(R) - \psi(z) = \int_z^R E dz \Rightarrow \psi(z) - \int_z^R E dz = -2\ell_1 \ln \frac{z}{R}$$

$$= -2\pi R^2 \ell_1 \ln \frac{z}{R}, \quad \ell_1 = \pi R^2 \epsilon P$$

$$2) a) z \geq R \quad E = \frac{2\ell_1}{z}, \quad \ell_1 = \int_R^R 2\pi z P(z) dz = 2\pi a \int_R^R z^2 dz =$$

$$= 2\pi a \frac{R^3}{3} \Rightarrow E = \frac{4\pi a R^3}{3z}$$

$$-\psi(z) + \psi(R) = - \int_2^R \frac{4\pi a R^3}{3z} dz; \quad \psi(z) = -\frac{4\pi a R^3}{3} \ln \frac{z}{R}$$

б)  $z \leq R$

$$2\pi r^2 D = 4\pi \int_0^z 2\pi r^2 \rho(z) dz$$

$$2D = 4\pi \int_0^z a z^2 dz = \frac{4\pi}{3} a z^3$$

$$D = \frac{4\pi}{3} a z^2 \quad E = \frac{D}{E} = \frac{4\pi}{3\epsilon} a z^2$$

$$\psi(z) - \psi(R) = \int_R^z \frac{4\pi}{3\epsilon} a z^2 dz = \frac{4\pi}{9\epsilon} a (R^3 - z^3)$$

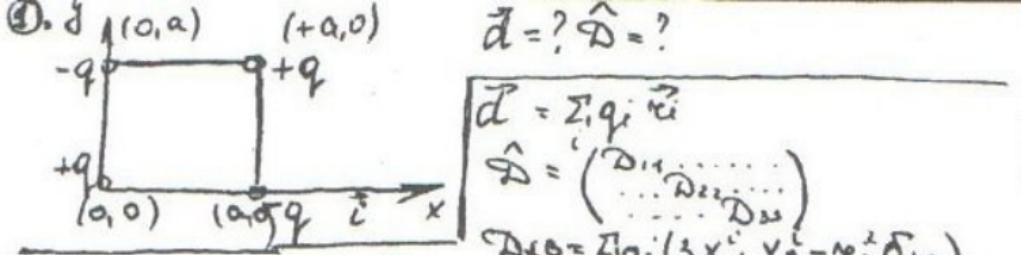
$$j = a z^2$$

$$\oint H d\vec{r} = 4\pi J$$

$$2\pi r^2 H = 4\pi \int_0^z j 2\pi r^2 dz = 4\pi \int_0^z a z^2 2\pi r^2 dz$$

$$2H = \frac{4\pi a z^3}{9} \quad H = \frac{4\pi a z^2}{9}$$





$$D_{ij} = D_{ji}$$

$$x_1 \rightarrow x$$

$$x_2 \rightarrow y$$

$$x_3 \rightarrow z$$

$$\textcircled{1} \quad \vec{d} = q \cdot 0 - q \cdot \vec{j} + q (\vec{a}_i + \vec{a}_j) + q \vec{a}_{ij} = 0$$

$$\textcircled{2} \quad D_{11} = D_{22} = \sum_i q_i (3(x^i)^2 - D_{11} u_i^2) =$$

$$= +q(3 \cdot 0 - 0) - q(3 \cdot 0 - a^2) +$$

$$+ q(3a^2 - 2a^2) = q(3a^2 - a^2) = 0$$

$$D_{12} = \sum_i q_i (3x^i y^i) = 3qa^2;$$

$$D_{21} = D_{12} = 3qa^2;$$

$$D_{13} = \sum_i q_i (3x^i z^i) = 0; \quad D_{31} = D_{13} = 0$$

$$D_{22} = \sum_i q_i (3(y^i)^2 - u_i^2) = q(0) - q(3a^2 - a^2) +$$

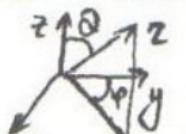
$$+ q(3a^2 - 2a^2) - q(3 \cdot 0 - a^2) = 0$$

$$D_{23} = \sum_i q_i (3y^i z^i) = 0 = D_{32}$$

$$D_{33} = q_i (3z^i - u_i^2) = -\sum_i q_i u_i^2 = -[q \cdot 0 - q \cdot a^2 + q \cdot a^2 -$$

$$\hat{D} = \begin{pmatrix} 0 & 3qa^2 & 0 \\ 3qa^2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad -q \cdot a^2] = 0;$$

$$\frac{dJ}{d\Omega} = \frac{\hat{D} \times \vec{n}}{144\pi c^2}, \quad |\vec{n}| = 1$$



$$\vec{n} = \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix}$$

$$\text{Hexagon } \frac{dJ}{d\Omega} = f(t); \quad \langle \frac{dJ}{d\Omega} \rangle_t = \lim_{n \rightarrow \infty} \frac{1}{n} \int_0^{2\pi} \frac{dJ}{d\Omega} dt \neq f(t);$$

$$②. \quad \begin{array}{c} -q \\ \xleftarrow[b,w]{} \\ ② +2q \\ \xrightarrow[a]{} \\ ③ -q \end{array}$$

$\nabla q_i(x, y) = 0$

$z_1 = 0; \quad z_2 = +a + b \sin wt; \quad z_3 = -a + b \sin wt$

$D_{1p} = \sum_i q_i (3 \Delta_x x_p^i - \Delta_y x_p^i); \quad D_{11} = \sum_i q_i (3 x_i^2 - v_i^2) = - \sum_i q_i v_i^2 = -(2q \cdot 0 - q z_2^2 - q z_3^2) = q(z_2^2 + z_3^2) = q(a^2 + b^2 \sin^2 wt + 2ab \sin wt + a^2 + b^2 \sin^2 wt - 2ab \sin wt) = 2q(a^2 + b^2 \sin^2 wt)$

$$D_{12} = \sum_i q_i (3 x^i y^i) = 0 = D_{21}$$

$$D_{22} = \sum_i q_i (3 y_i^2 - v_i^2) = - \sum_i q_i v_i^2 = 2q(a^2 + b^2 \sin^2 wt)$$

$$D_{23} = \sum_i q_i (3 y^i z^i) = 0 = D_{32}$$

$$D_{31} = D_{13} = 0; \quad D_{33} = \sum_i q_i (3(z^i)^2 - v_i^2) = \sum_i 2q_i v_i^2 = 2 \sum_i q_i v_i^2 = -4q(a^2 + b^2 \sin^2 wt)$$

$$\hat{\mathcal{D}} = \begin{pmatrix} 2q(a^2 + b^2 \sin^2 wt) & 0 & 0 \\ 0 & 2q(a^2 + b^2 \sin^2 wt) & 0 \\ 0 & 0 & -4q(a^2 + b^2 \sin^2 wt) \end{pmatrix}$$

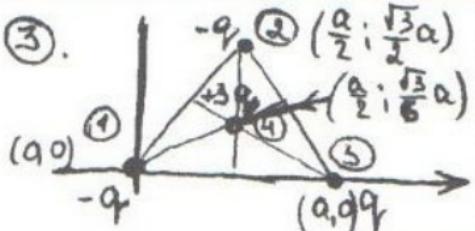
$$\hat{\mathcal{D}} = \begin{pmatrix} qb^2 w \sin wt \cos wt & 0 & 0 \\ 0 & -w & 0 \\ 0 & 0 & -4q w b^2 \sin^2 wt \cos wt \end{pmatrix}$$

$$\hat{\mathcal{D}} = \dots \quad \frac{d\hat{\mathcal{D}}}{ds} = \frac{[\hat{\mathcal{D}} \times \vec{n}]^2 \vec{n}}{144\pi c^5}; \quad \hat{\mathcal{D}} \times \vec{n} \quad \textcircled{E}$$

$$\textcircled{E} \quad \begin{pmatrix} -8q w^2 b^2 \sin^2 wt & 0 & 0 \\ 0 & -8q w^3 b^2 \sin^2 wt & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \begin{pmatrix} \text{cos}\theta \\ \text{sin}\theta \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} -8q w^3 b^2 \sin^2 wt \cos\theta \\ -8q w^3 b^2 \sin^2 wt \sin\theta \\ 0 \end{pmatrix} \quad \frac{d\hat{\mathcal{D}}}{ds} = \frac{64 q^2 w^6 b^4 \sin^2 wt}{144\pi c^5}$$

③.



$$\text{Q. } \vec{a} = \sum_i q_i \vec{u}_i =$$

$$= -\vec{q}_1 \cdot 0 - \vec{q}_1 \left( \frac{a}{2} \vec{i} + \frac{\sqrt{3}}{2} \vec{j} \right) -$$

$$- \vec{q}_1 (\vec{a} \vec{i}) + 3 \vec{q}_1 \left( \frac{a}{2} \vec{i} + \frac{\sqrt{3}}{2} \vec{a} \vec{j} \right) =$$

$$= \vec{i} \left( -\frac{9a}{2} + \frac{5\sqrt{3}a}{6} \right) \oplus$$

$$\textcircled{+} \quad \vec{j} \left( -\frac{\sqrt{3}}{2} \vec{q}_1 a - \vec{q}_1 a + \frac{\sqrt{3}}{2} \vec{q}_1 a \right) = \vec{q}_1 \vec{a} \vec{i} - \vec{q}_1 \vec{a} \vec{j} =$$

$$= \vec{q}_1 a (\vec{i} - \vec{j}).$$

$$\textcircled{0} \quad D_{11} = \sum_i q_i (3(x_i^2) - v_i) = -\vec{q}_1 (3 \cdot 0 - 0) -$$

$$-\vec{q}_1 (3 \frac{a}{2} \cdot \frac{a}{2} - a^2) - \vec{q}_1 (3 \cdot a^2 - a^2) + 3 \vec{q}_1 \left( \frac{a^2}{4} - \frac{a^2}{3} \right) =$$

$$= 0 - \vec{q}_1 \left( 3 \frac{a^2}{4} - a^2 \right) - \vec{q}_1 \cdot 2a^2 = -\vec{q}_1 \frac{a^2}{2};$$

$$D_{12} = \sum_i q_i \left( 3x_i y_i - \frac{v_i}{z_i} \right) = -\vec{q}_1 \cdot 0 - \vec{q}_1 \left( 3 \frac{a}{2} \cdot \frac{\sqrt{3}}{2} a \right) +$$

$$+ \vec{q}_1 a + 3 \vec{q}_1 \left( \frac{a}{2} \cdot \frac{\sqrt{3}}{6} a \right) - \vec{q}_1 a = 0 \Rightarrow D_{12} = 0$$

$$D_{13} = \sum_i q_i (3x_i z_i - v_i) = 0 \Rightarrow D_{13} = 0$$

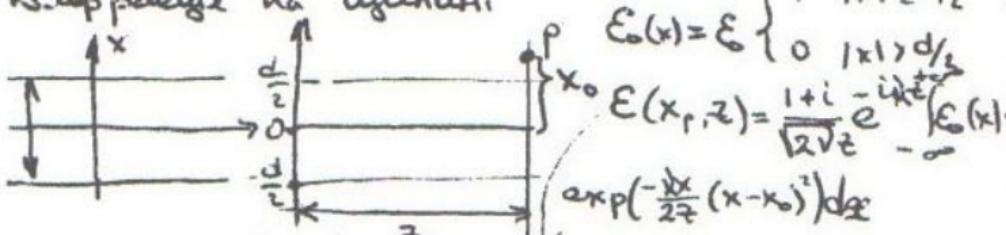
$$D_{22} = \sum_i q_i (3y_i^2 - v_i^2) = -\vec{q}_1 (0) - \vec{q}_1 \left( 3 \left( \frac{\sqrt{3}}{2} a \right)^2 - a^2 \right) -$$

$$- \vec{q}_1 (0 - a^2) + 3 \vec{q}_1 \left( 3 \left( \frac{\sqrt{3}}{6} \right)^2 - \left( \frac{a}{\sqrt{5}} \right)^2 \right) = -\frac{9a^2}{2};$$

$$D_{23} = D_{32} = 0;$$

$$D_{33} = - \sum_i q_i v_i^2 = - \left( 0 - \vec{q}_1 a^2 - \vec{q}_1 a^2 + 3 \vec{q}_1 \cdot \frac{a^2}{3} \right) =$$

$$= \vec{q}_1 a^2;$$



$$\begin{aligned} \xi(x) &= \xi \left\{ \begin{array}{ll} 0 & |x| > d/2 \\ 1 & |x| \leq d/2 \end{array} \right. \\ \xi(x_p, z) &= \frac{1+i}{\sqrt{2\pi z}} e^{-\frac{i k z}{2}} \int_{-\infty}^{\infty} \xi(x) \exp\left(-\frac{ikx}{2z}(x-x_p)^2\right) dx \\ &\quad \times \exp\left(-\frac{ikx}{2z}(x-x_0)^2\right) dx \end{aligned}$$

$$\xi(x_p, z) = \xi_0 \frac{1+i}{\sqrt{2\pi z}} e^{-\frac{i k z}{2}} \int_{-\infty}^{\infty} \xi(x) \exp\left(-\frac{ikx}{2z}(x-x_0)^2\right) dx$$

$$\therefore (x_0, z) = \xi_0 \frac{1+i}{\sqrt{2\pi z}} e^{-\frac{i k z}{2}} J_1(x_0)$$

$$J_1(x_0) = \int_{-d/2}^{d/2} \exp\left(-\frac{ikx}{2z}(x-x_0)^2\right) dx, z = x - x_0 \sqrt{\frac{K}{\pi z}}$$

$$J_1(x_0) = \sqrt{\frac{\pi z}{K}} \int_{-\zeta_2}^{\zeta_2} \exp\left(-\frac{ik\zeta}{2}\right) d\zeta.$$

$$\zeta_1 = -\left(\frac{d}{2} + x_0\right) \sqrt{\frac{K}{\pi z}}, \quad \zeta_2 = +\left(\frac{d}{2} - x_0\right) \sqrt{\frac{K}{\pi z}}$$

$$C(\alpha) = \int \cos \frac{ikt}{2} dt \quad S(k) = \int \sin \frac{ikt}{2} dt$$

$x = 0$  интервал  $\Phi$  времени

$$\ell^2(\alpha_1, \alpha_2) = (C(\alpha_1) - C(\alpha_2))^2 + (S(\alpha_1) - S(\alpha_2))^2$$

$$J(x_0, z) = \frac{1}{2} J_0 \ell^2(\zeta_1, \zeta_2), \quad J_0 - \text{интервал наведения}$$

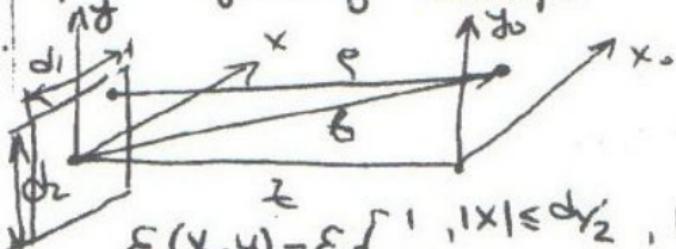
$$\frac{J(x_0)}{J_0} = \frac{1}{2} \ell^2(\zeta_1, \zeta_2)$$

$$\zeta_1 = -\alpha(1+\rho), \quad \zeta_2 = \alpha(1-\rho)$$

$$\alpha = \sqrt{\frac{Kd^2}{4\pi z}} = \sqrt{2N_f} \quad N_f - \text{число} \Phi \text{ времени}$$

$$\rho = 2k_0 l + 1$$

шараханде Траунгофера на кручениш  
принескуюшы отвори.



$$E_0(x, y) = E_0 \begin{cases} 1, & |x| \leq d_1/2, |y| \leq d_2/2 \\ 0, & \text{зәмбән} \end{cases}$$

$$E_0(k_x, k_y) = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} E_0(x, y) e^{i(k_x x + k_y y)} dx dy$$

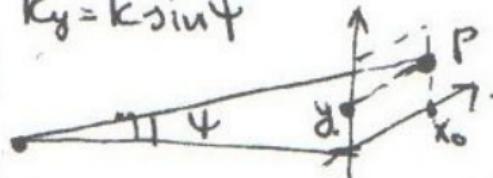
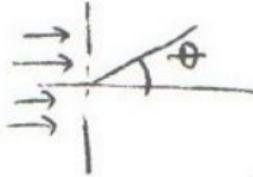
$$E_0(k_x, k_y) = E_0 \int_{-\frac{d_1}{2}}^{\frac{d_1}{2}} dx \int_{-\frac{d_2}{2}}^{\frac{d_2}{2}} dy e^{i(k_x x + k_y y)} = E_0 d_1 d_2 \operatorname{sinc}\left(\frac{k_x d_1}{2}\right) \operatorname{sinc}\left(\frac{k_y d_2}{2}\right)$$

$$J(P) = \frac{c}{8\pi} \frac{1}{(KB)^2} S_0(k_x, k_y) - \text{просторова спектралык нусқаны}$$

$$S_0(k_x, k_y) = (E_0(k_x, k_y))^2$$

$$k_x = k \sin \theta$$

$$k_y = k \sin \psi$$



$$J(\theta, \psi) = J_{\max} \operatorname{sinc}^2\left(\frac{td_1^2 \sin \theta}{\lambda}\right) \cdot \operatorname{sinc}^2\left(\frac{nd_2 \sin \psi}{\lambda}\right)$$

$$J_{\max} = J_0 \left( \frac{d_1 d_2}{KB} \right)$$

Дисперсия Гауссова пучка

$$E_0(r) = E_0 \exp\left(-\frac{r^2}{2\sigma_r^2}\right),$$

$$E(k_z) = E_0 \cdot 2\pi \exp\left(-\frac{k_z^2 \rho_0^2}{2}\right)$$

$$Q \ll 1 \quad k_z = k_Q$$

$$J(Q) = J_{\max} \exp\left[-\frac{2\pi Q \rho^2}{a}\right]$$

$$J_{\max} = J_0 \left(\frac{2\pi \rho}{a}\right)^2$$

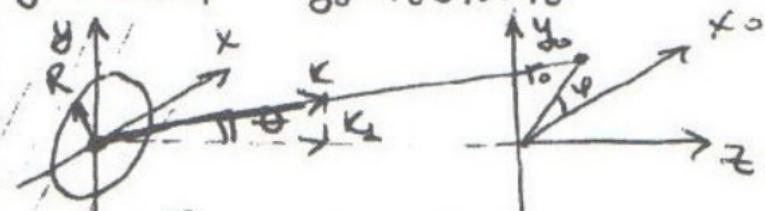
$$J_0 = \frac{c}{8\pi} (E_0)^2,$$

$\rho_0$  - начальный радиус пучка

Дифракция на круглом отверстии

$$x = r \cos \varphi \quad x_0 = r_0 \cos \varphi_0$$

$$y = r \sin \varphi \quad y_0 = r_0 \sin \varphi_0$$



$$\sin \theta = \frac{r_0}{R}, \quad K \sin \theta = K_1$$

$$K_x = K \frac{x_0}{R} = K_1 \cos \varphi_0; \quad K_x x + K_y y = K_1 r \cos(\varphi - \varphi_0)$$

$$K_y = K \frac{y_0}{R} = K_1 \sin \varphi_0; \quad \text{2nd}$$

$$\sum(K_1, \varphi_0) = \int_0^{\infty} \sum_0(r) dr \cdot \int_{-\pi}^{\pi} \exp(i K_1 r \cos(\varphi - \varphi_0)) d\varphi$$

$$\int_{-\pi}^{\pi} \exp(i \theta \cos(\varphi - \varphi_0)) d\varphi = 2\pi, (\alpha)$$

• Г-о-р-я бесконечного породы  
переворотение

$$\sum(K_1) = 2\pi \int_0^{\infty} \sum_0(r) J_0(K_1 r) r dr - \text{Фурье-бесконечное}\braket{хенкеле} \text{нечеткого порода}\braket{круглого отверстия} R:$$

$$\sum_0(r) = \sum_0(0), r > R, \quad \sum(K_1) = \sum_0 \cdot 2\pi \int_0^{\infty} J_0(K_1 r) r dr =$$

$$= \sum_0 \cdot 2\pi R^2 \frac{J_1(K_1 R)}{K_1 R}, \quad \text{где вычисляется.}$$

$$\int x \cdot J_0(x) dx = x J_1(x), \quad \theta \ll 1, \quad \sin \theta = \theta$$

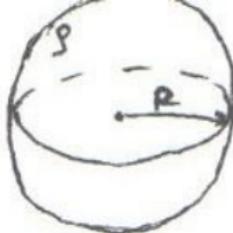
$$K_1 = K \sin \theta = K \theta$$

$$J(P) = \frac{C}{8\pi} \frac{1}{(AB)^2} S_0(K_x, K_y)$$

$$S_0(K_x, K_y) = [\sum_0(K_x, K_y)]^2 \Rightarrow J(\theta) = J_{\max} \left( \frac{J_1(2\pi \theta R / K)}{\pi \theta R / K} \right)^2$$

$J_{\max} = J_0 \left( \frac{\pi R^2}{AB} \right)^2$ ,  $J_0$  - интеграл падающей хвани  
половине куговой ширине центрального  
максимума диф. картине

$$\Delta \theta = 1.22 \cdot \lambda / R$$



$$E = ?$$

1)  $z < R$

$$\oint \vec{E} d\vec{s} = 4\pi Q$$

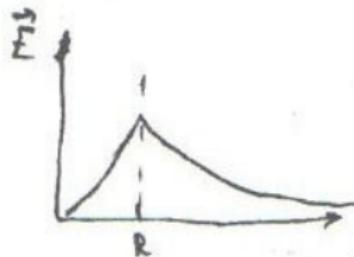
$$E \cdot 4\pi z^2 = 4\pi \rho \frac{4}{3}\pi z^3$$

$$E = \frac{4}{3}\pi \rho z^2$$

2)  $z > R$

$$E \cdot 4\pi R^2 = 4\pi \rho \frac{4}{3}\pi R^3$$

$$E = \frac{4}{3}\rho \frac{R^3}{z^2}$$



$$\varphi_2 = - \int E dz = - \frac{2}{3}\pi \rho z^2 + C_0 \quad (z \leq R)$$

$$\varphi = \frac{4}{3}\rho R^3 z^{\frac{1}{2}} + C \quad (z > R)$$

"o( $\varphi(\infty) = 0$ )"

$$- \frac{2}{3}\pi \rho R^2 z + C_0 = \frac{4}{3}\rho R^2 \pi$$

$$C_0 = 2\pi R^2 \rho$$

$$\varphi = - \frac{2}{3}\pi \rho z^2 + 2\pi R^2 \rho$$

іднр. джерело в більшій зоні. Внаслідок  
 $R_H = \left| \frac{E_0}{H_{\alpha^0}} \right| \cdot$  діє гарячий ефект.

Поглядання:

$$\psi(\vec{r}, t) = \frac{q\varepsilon}{r} + \frac{\vec{n} \cdot \vec{p}(t')}{r^{1/2}} + \frac{\vec{n} \cdot \vec{p}(t')}{c|r|} - \text{Вих. висок}\newline
\text{частотне}\newline
\text{збурення}\newline
\text{поміжний}\newline
\text{випадок}$$

$$q\varepsilon = \int p d^3 r'$$

Спершу вибираємо координати

$$\vec{E} \cdot \vec{p} = P_0 \exp(-i\omega t') \cos\theta$$

$$\vec{H} = \vec{e}_z \cdot H_{\alpha^0} \exp(-i\omega t')$$

$$\vec{E} = (\vec{e}_z \cdot E_z + \vec{e}_{\theta} \cdot E_{\theta}) e^{-i\omega t'}$$

$$H_{\alpha^0} = (1 + i \frac{\lambda}{2\pi|r|}) P_0 \sin\theta$$

$$E_{\theta} = \left[ 1 - \frac{\lambda}{2\pi|r|} \right]^2 + i \frac{\lambda}{2\pi|r|} P_0 \sin\theta$$

$$E_z = L - 2i \left( \frac{\lambda}{2\pi|r|} \right) \left( 1 + i \frac{\lambda}{2\pi|r|} \right) P_0 \cos\theta$$

$$P_0 = -\frac{\omega}{c} \cdot \frac{P_0}{r^{1/2}}$$

$\theta = \Theta$  — в більшій зоні відсутнє навколо

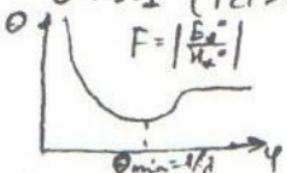
$$\left| \frac{E_{\theta}}{H_{\alpha^0}} \right| = \left| \frac{1 - \left( \frac{\lambda}{\Theta} \right)^2 + i \left( \frac{\lambda}{\Theta} \right)}{1 + i \frac{\lambda}{\Theta}} \right| = \frac{1}{\Theta} \sqrt{\frac{1 - \Theta^2 + \Theta^2}{1 + \Theta^2}}$$

$$\Theta = \frac{\lambda r}{2\pi}$$

$$\Theta = \frac{\pi}{2} \Rightarrow E_z = 0$$

$$\Theta \gg 1 \quad (|r| \gg \lambda) : \left| \frac{E_{\theta}}{H_{\alpha^0}} \right| = 1$$

$$F = \left| \frac{E_{\theta}}{H_{\alpha^0}} \right|$$

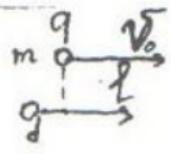


в більшій зоні:

$$\left| \frac{E_{\theta}}{H_{\alpha^0}} \right| \gg 1$$

$$F = \frac{1}{\Theta}$$

$$\frac{dF}{d\Theta} = 0$$



$$m \ddot{x} = q E$$

$$m \frac{dV}{dt} = q E$$

$$E = -q \text{grad}$$

$$E = \vec{p} \left\{ \frac{\vec{e}_z \cdot \vec{r}^3 - 3x\vec{r}^2}{r^5}; \quad ; \quad \right\}$$

$$\frac{dV}{dt} = \frac{q E}{m}$$

$$V - V_0 = -\frac{q E}{m} t$$

$$V = \frac{q E}{m} t + V_0$$

$$\dot{J} = n q V$$

$$H = \frac{1}{c} \int \frac{J dV}{2} = \frac{n q}{c} \int \frac{V dV}{2} dx dy$$

$$E = \frac{1}{c} [ [\vec{A} \times \vec{n}] \times \vec{n} ] = \omega \vec{t} + \vec{A} \times \vec{n}$$

Бард 9 , то совместное с  $W_0$ , норотиво  
акустичус -  $A_0$ , відбувається динамічне  
виродження. Іншими доказами єдиним  
акустичус коямбанс



$$W = \frac{m \omega^2 R^2}{2}$$

$$P_{\text{акуст.}} = \frac{dW}{dt}$$

$$P_{\text{норот. вир.}} = J = \frac{2}{5} \frac{\dot{P}}{c^2}$$

Перебреже фасрастареп гіасектр синанс  $E_0$ ,  
уго білбіи більши за  $E_\infty$  Значиме  
згноду швидкісті  $e/c$ и хвилі на зачесоні,  
уго білбіи менше за резонансу.

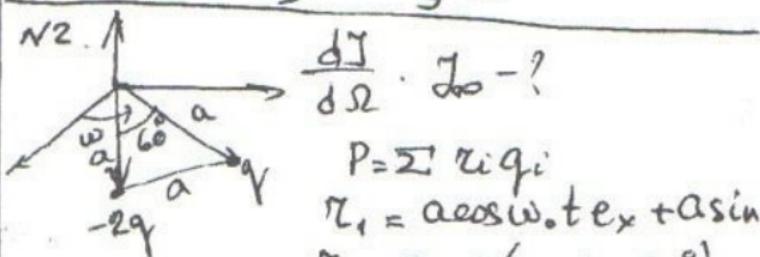
$$P-\text{нл}: \frac{E_0 - E_\infty}{\omega_0^2 - \omega^2}$$

$$E = E_\infty + \frac{(E_0 - E_\infty)}{\omega_0^2 - \omega^2} \omega^2$$

$$V_{zp} = \frac{dw}{dk}; k(\omega) = \frac{\omega}{c} \sqrt{\frac{E}{\omega}}$$

$$V_{zp} = \frac{dw}{d(\frac{\omega}{c} \sqrt{E(\omega)})}$$

$$V_{zp} = \frac{g_c - \epsilon_0 (1 - \frac{4}{3} \epsilon_0)}{\frac{13}{9} \epsilon - \frac{4}{3} \epsilon_0^2}$$



$$P = \sum r_i q_i$$

$$r_1 = a \cos \omega_0 t e_x + a \sin \omega_0 t e_y$$

$$r_2 = a \cos (\omega_0 t + 60^\circ) e_x + \dots$$

$$E = E_0 \exp \left\{ -\frac{t}{T} + \vec{k} \vec{r} \right\} \quad t > 0$$

$$E(\omega) = ?$$

$$E(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{-i\omega t} dt$$

Всегда  $+q$  рухається прямолінійно і  
має в нормальній сис. що  $V_0$ . На здог  
гие гальмівна сис., яка пропорційна  
до розрізняння. Задане вектор-параметр  
зупинювання чогось залежу в даній зоні.  
Розв'язання:



$$F_2 = -\frac{dV^2}{dt} = -dV_z \vec{e}_z, \quad d > 0$$

2ий закон Ньютона

$$m \ddot{\vec{z}} = \vec{F}_2.$$

$$m \ddot{\vec{z}} = -dV_z \vec{e}_z = -d \dot{\vec{z}}$$

$$m \ddot{\vec{z}} + d \dot{\vec{z}} = 0 \quad \ddot{z} = k$$

$$m \ddot{k} + dk = 0,$$

$$m \frac{dk}{dt} + dk = 0; \quad \frac{dk}{k} = -\frac{d}{m}, \quad \ln \frac{k}{k_0} = -\frac{d}{m} t,$$

$$k = k_0 e^{-\frac{d}{m} t}; \quad z = \dot{z} = \dot{z}_0 e^{-\frac{d}{m} t};$$

$$v(t=0) = v_0 = \dot{z}_0, \quad v = v_0 e^{-\frac{d}{m} t}$$

$$A(\vec{r}, t) = \frac{1}{c(\vec{r})} \quad q \quad v(t') \vec{e}_z; \quad t' = t - \frac{|\vec{r}|}{c}$$

$$A(\vec{r}, t) = \frac{1}{c} \quad -\frac{d}{m} t - \frac{|\vec{r}|}{c} \rightarrow$$

# Бівім 49?

① По квадратній амплітуді миє  
спрощ.

$$J = \frac{at}{b^2 f^2} \quad \begin{matrix} \text{Знайдене} \\ \text{випадок} \end{matrix} \rightarrow \text{спектр} \\ \text{Dob. Page 21}$$

②  $\rho = r^2 e^{-\frac{r}{a}}$

Заряд розподілений у просторі з густинкою  
 $\rho$ . Знайдіти волу ми номенклатури заряду  $Q$ ,

що є відомою  $Q = A = \text{const.}$

Розв'язання: Нехай  $\rho = Ar^2 e^{-\frac{r}{a}}$

Ф-я Оміора - Тейса:

$$\oint E dS = 4\pi Q, \quad Q = \int_0^R 4\pi A r^4 e^{-\frac{r}{a}} dr;$$

$$\oint E dS = 4\pi \int_0^R 4\pi A r^4 e^{-\frac{r}{a}} dr$$

$$E = 4\pi R^2 = \frac{96\pi^2 A}{5} \int_0^R r^4 e^{-\frac{r}{a}} dr,$$

$$E = \frac{4\pi A}{R^2} \int_0^R r^4 e^{-\frac{r}{a}} dr$$

Знайдемо  $\text{const } A$ :

$$Q = \int_0^\infty r^4 e^{-\frac{r}{a}} dr = 4\pi A \frac{4!}{a^5}$$

$$\int_0^\infty r^n e^{-\frac{r}{a}} dr = \frac{n!}{a^{n+1}}$$

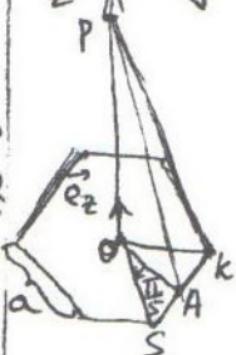
$$A = \frac{Q}{3\pi}$$

$$E = \frac{4Q}{3R^2} \int_0^R r^4 e^{-\frac{r}{a}} dr$$

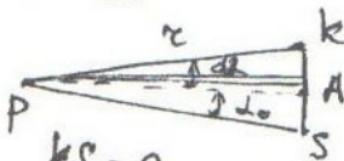
$$Q = \int_0^R E(r) dr$$

2) Тірабиньній нелінійний з міцінною  
системою  $\rho$ , сторона  $a$ . Значення на  
інтенсивні на осі симетрії системи  
 $E(0, 0, 2)$ ,  $\varphi(0, 0, 2)$

Погляд згорі!



$$\rho = \frac{dq}{dl}$$



$$dq = \rho dl = \rho r dd$$

$$d\varphi = \frac{dq}{r^2} = \rho d\alpha, \alpha \in [-\alpha_0, \alpha_0]$$

$$\tan \alpha_0 = \frac{AS}{AP} = \frac{\alpha/2}{\sqrt{z^2 + OA^2}}; OA = \frac{a}{2} \tan \frac{\pi}{S} = d$$

$$\tan \frac{\pi}{S} = \frac{OA}{AS} = \frac{OA}{\alpha/2} \rightarrow OA = \frac{a \tan \frac{\pi}{S}}{2}$$

$$\varphi_1 = \int_{-\alpha_0}^{\alpha_0} d\varphi = \rho \int_{-\alpha_0}^{\alpha_0} dd = \rho 2 \alpha_0 = 2 \rho \arctan \frac{a}{\sqrt{z^2 + d^2}}$$

$$\varphi(0, 0, z) = S(\varphi_1)$$

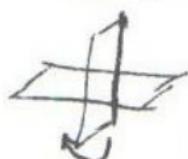
$$E = -\operatorname{grad} \varphi = -\frac{\partial \varphi}{\partial z} \vec{e}_z$$

Задача №55

Знайдіть власні матриці гідрав. прописки.  
 якої будуть ненульовими лише в крестах, які  
 є симетрією крестика  $C_3$  і не симетрії

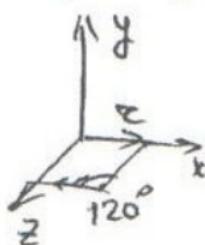
Розв'язання:

$$C_3 \rightarrow \varphi = \frac{2\pi}{3} = 120^\circ$$



$$\varepsilon_{ij} = \begin{pmatrix} \varepsilon_{xx} & 0 & 0 \\ 0 & \varepsilon_{yy} & 0 \\ 0 & 0 & \varepsilon_{zz} \end{pmatrix}$$

$$\varepsilon'_{i'j'} = a_{i'i} \varepsilon_{ij} a_{ij'}$$



$$\begin{pmatrix} -\frac{1}{2} \\ -\frac{\sqrt{3}}{2} \\ 0 \end{pmatrix} = \begin{pmatrix} \dots \\ \dots \\ \dots \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$a_{i'i} = \begin{pmatrix} -\frac{1}{2} & \dots \\ -\frac{\sqrt{3}}{2} & \dots \\ 0 & \dots \end{pmatrix}$$

$$\begin{pmatrix} \frac{\sqrt{3}}{2} \\ -\frac{1}{2} \\ 0 \end{pmatrix} = \begin{pmatrix} \dots & \frac{\sqrt{3}}{2} \\ \dots & -\frac{1}{2} \\ \dots & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}; \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} \dots & 0 \\ \dots & 0 \\ \dots & 1 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$a_{i'i} = \begin{pmatrix} -\frac{1}{2} & \frac{\sqrt{3}}{2} & 0 \\ -\frac{\sqrt{3}}{2} & -\frac{1}{2} & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad \varepsilon'_{i'j'} = \begin{pmatrix} -\frac{1}{2} & \frac{\sqrt{3}}{2} & 0 \\ -\frac{\sqrt{3}}{2} & -\frac{1}{2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \varepsilon_{xx} & 0 & 0 \\ 0 & \varepsilon_{yy} & 0 \\ 0 & 0 & \varepsilon_{zz} \end{pmatrix} =$$

$$= \varepsilon'_{i'j'} = \begin{pmatrix} \dots & \dots & 0 \\ \dots & \dots & 0 \\ \dots & \dots & 0 \end{pmatrix} = \begin{pmatrix} \varepsilon'_{11} & \varepsilon'_{12} & 0 \\ \varepsilon'_{21} & \varepsilon'_{22} & 0 \\ 0 & 0 & \varepsilon'_{33} \end{pmatrix}$$

$$\varepsilon_{sym} = \delta_B \varepsilon'_{ij''}$$

$$\delta_n = \begin{pmatrix} 1 & 0 \\ 0 & 1 \\ 0 & -1 \end{pmatrix} \quad \delta_B = \begin{pmatrix} -1 & 0 \\ 0 & -1 \\ 1 \end{pmatrix}$$

2)  $K_{fac}$

### Більш 50

Дано коло R, поділено на 3 рівні сектори. З них 2 зареджені позитивно - негативно, 1 чисто зареджені позитивно +B, негативно .- " + 2B  
 $d = ? \quad r = ?$



$$\text{Розв'язання: } \varphi = \frac{2\pi}{3}$$

$$P = \sum P_i, \quad dl = R d\varphi$$

$$x = R_0 \cos \varphi, \quad y = R_0 \sin \varphi$$

$$dP = \rho \times dl, \quad \text{де } \rho - \text{плотність} \quad \text{зусімка}$$

$$P = \int_{-\frac{\pi}{3}}^{\frac{\pi}{3}} \rho R_0^2 d\varphi \cos \varphi = \rho R_0^2 \sin \varphi \Big|_{-\frac{\pi}{3}}^{\frac{\pi}{3}} = \sqrt{3} \rho R_0^2$$

$$\vec{P}_1 = \sqrt{3} B R_0^2, \quad \vec{P}_2 = \sqrt{3} B R_0^2, \quad \vec{P}_3 = 2\sqrt{3} B R_0^2$$

$$P = \sum P$$

Оцінити час обертання на один.

$$\frac{m \omega^2}{2} = \frac{l_0^2}{r^2} \quad \omega_0 = \sqrt{\frac{l_0^2}{mr_0}}$$

$$T = \frac{2\pi \omega_0}{\nu} \approx 10^{-16} \text{ c}$$

За час обертання  $\approx 10^6$  обертів

2) Квадратна симетрия з є симетрією  $\alpha$ ,  
 як і не має відмінного напруження.  
 маска хвиль. Задані  $E_{\text{наг}}$ ,  $x_{\text{бок}}$ , яко  
 нагадує нам кутове  $Q$ .  
 Розв'язок:

$$H = H_0 e^{i(\omega t - \vec{k} \cdot \vec{r})}$$

$$E_{\text{наг}} = -\frac{1}{c} \frac{\partial \Phi}{\partial t}$$

$$\operatorname{Re}(\Phi) = \iint_S H dS = \iint_{\substack{y_1 \leq y \leq y_2 \\ -x_2 \leq x \leq x_2}} H_0 e^{i(\omega t - \vec{k} \cdot \vec{r})} dx dy =$$

$$= \iint \left| H \right| \cos Q \cos (\omega t - \vec{k} \cdot \vec{r}) = \iint \left| H \right| \cos Q \times$$

$$\times \cos (\omega t - (k_x x + k_y y + k_z z)) =$$

$$= - \iint \left| H \right| \cos Q \cos (\omega t + k \cos Q) dx dy =$$

$$= - \left| H \right| \cos Q (\omega t + k \cos Q) \cdot Q^2$$

$$\Phi_r = w \cdot \Phi$$

$$E_{\text{наг}} = \frac{w}{c} \Phi (\sin \omega t + k \cos Q) \cos Q$$

Бүгүнгөмөн радиоактивтілік ерекшіліктер

$$f_{\text{күн}} \gg f_{n.r.}$$

$$\ddot{P}_S = \frac{2 l_0^2}{3 C^3} < (\ddot{\tau}) >$$

Мысалы нәсекомые  $|f_{n.r.}| \gg 0$

$$f = f_{\text{күн}} + f_{n.r.}$$

$$f_{n.r.} = \frac{2 l_0^2}{3 C^3} \ddot{\tau}, \quad \ddot{\tau} = - \frac{l_0^2}{m(r_0)^3} \ddot{\tau}_0 < \frac{1}{x} > \neq \frac{1}{x}$$

$$J = P_S = \frac{2 l_0^2}{3 C^3 m^2} < \left| \frac{1}{\tau_0(t)} \right|^4 > = \frac{2 l_0^6}{3 C^3 m^2} \frac{1}{\tau_0^4}$$

$$E_{\text{күн}} = - \frac{l_0^2}{|r|}; \quad P_S \approx \frac{2 E_{\text{күн}}^4}{3 C^3 m^2 l_0^2}$$

$$E = T + U = \frac{m \dot{r}^2}{2} - \frac{l_0^2}{2}; \quad \frac{m \dot{r}}{|r|} = \frac{l_0^2}{|r|^2};$$

$$E = \frac{1}{2} E_{\text{күн}}; \quad \frac{dE}{dt} = - P_S$$

$$\frac{dE_{\text{күн}}}{dt} + b E_{\text{күн}}^4 = 0, \quad b = \frac{4}{3 C^3 l_0^2 m^2}, \quad t=0$$

$$\tau_0(t=0) = a_0, \quad E_{\text{күн}}(t=0) = W_0 = - \frac{l_0^2}{a_0}$$

$$\int \frac{dE_{\text{күн}}}{E_{\text{күн}}^4} = -b \int dt$$

$$E_{\text{күн}}(t=T) = - \frac{l_0^2}{\tau_0(T)} = \infty, \quad \text{ерекшілік} \rightarrow \text{егемен}$$

$$\frac{1}{3} E_{\text{күн}} \Big|_{W_0}^{\infty} = bT$$



$$T = \frac{a_0^3 C^3 m^2}{4 l_0^4}, \quad a_0 \sim 1 \text{ \AA}, \quad T = 10^\circ \text{C}$$